

Dark Matters - The trouble  
with searching for rare events

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# Looking for Rare Events

- First of all, what is considered rare?
- Certainly anything that happens once in the span of a lifetime. Maybe once in a year?
- “You’ll know one when you see one ... if you see one”
- More formal:

binary dependent variables with dozens to thousands of times fewer ones (events, such as wars, vetoes, cases of political activism, or epidemiological infections) than zeros (“nonevents”) - Logistic Regression in Rare Events Data

# Dealing with Rare Events

- So what can we do? Seem to be three solutions:
  - Wait
  - Predict from other information
  - Interpret the results we have

# Rare Events

- Let's assume that anything that happens  $< 1$  time per month is rare
- Could be many things... meteor strikes, government upheavals, floods, birthdays...
- Can we do anything predictive with this class of events?

# Answer: Maybe?

- Discrete physical phenomena tend to follow a power law distribution
- Can be used with earthquake magnitude
- Gutenberg-Richter law

$$N = 10^{a-bM}$$

N - number of earthquakes

M - magnitude of the earthquake

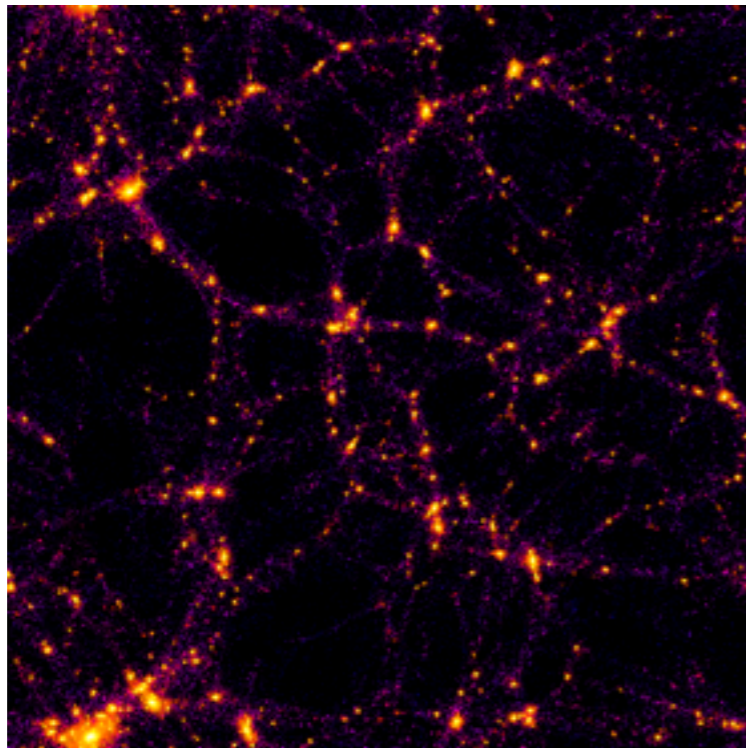
a,b - constants

# Earthquake Data

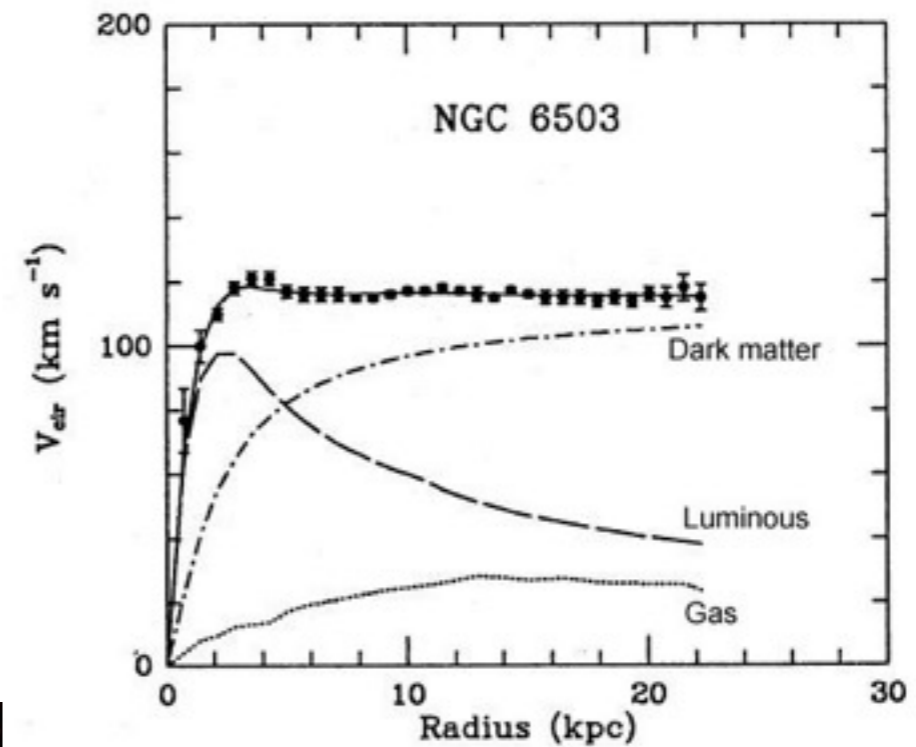
- Data on recent earthquakes is available here: <http://earthquake.usgs.gov/earthquakes/feed/v1.0/csv.php>
- Can you predict the number of occurrences of magnitude 5? 6?
- How accurate are the predictions?

# Dark Matter Introduction

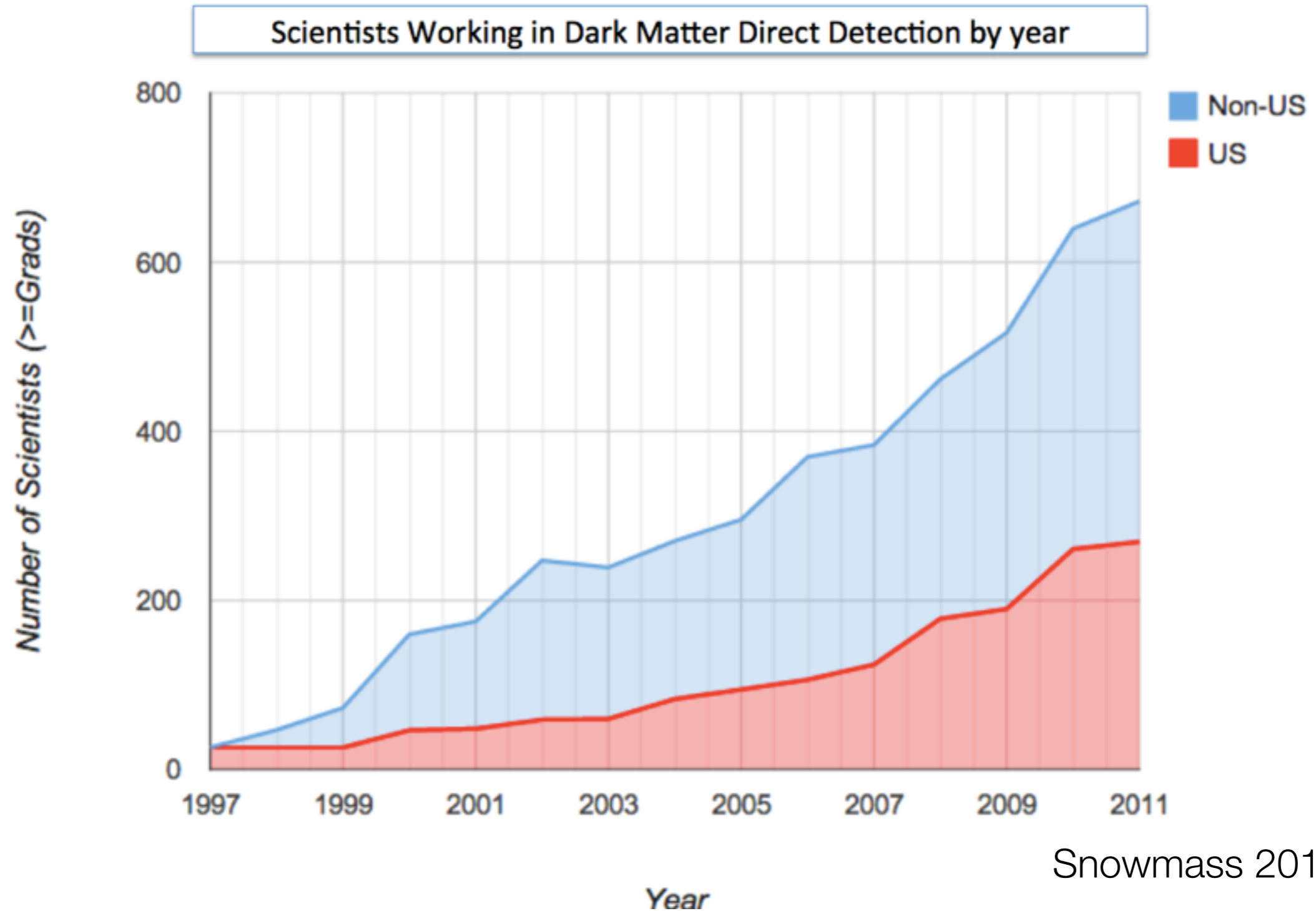
Structure formation with simulation



Separation of luminous and non-luminous mass



# A Popular Field...



Snowmass 2013



# Dark Matter

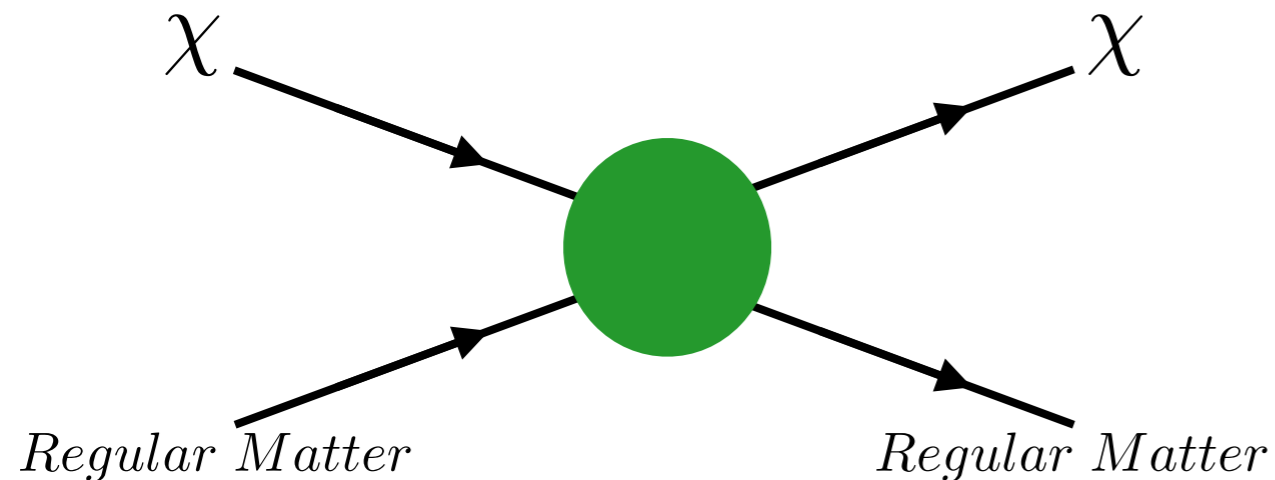
What do we know?

1. Long lived (survived until current day)
2. Non-baryonic (Hydrogen:Deuterium Ratio)
3. No EM interactions (haven't seen it)
4. 80% of all matter (rotation curves, CMB)
5. Non-relativistic (galactic structure formation, rotation curves)

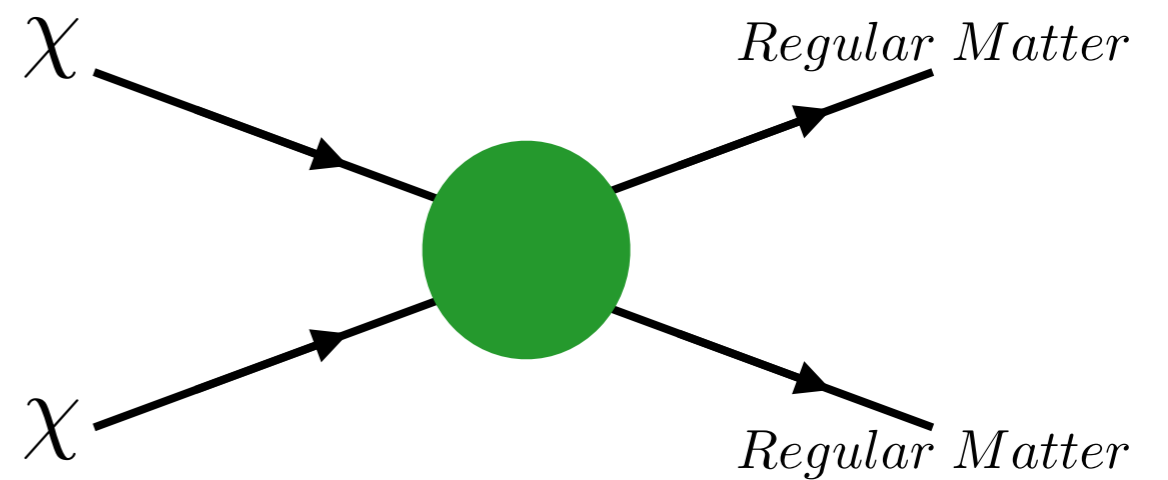
# Dark Matter Detection

Two primary types: direct and indirect

Interaction of particle  
inside the detector



Products of self-interaction  
outside the detector



# Dark Matter Math

$$\frac{dR}{dE_\nu} = \underbrace{\sigma_p}_{\text{cross-section}} \sum_A \underbrace{f(A)}_{\text{mass fraction of target}} \times \underbrace{S(A, E_R)}_{\text{phase space distribution of trapped DM}} \times \underbrace{I(A)}_{\text{coupling enhancement}} \times \underbrace{F^2(A, E_R)}_{\text{nuclear form factor}} \times \underbrace{g(A)}_{\text{nuclear recoil quenching factor}} \times \underbrace{\epsilon(E_\nu)}_{\text{detector response function}}$$

A - Atomic Number

$E_R$  - nuclear recoil energy

# Direct Detection Methods

- All experiments search for **small** energy deposits
- Make some assumptions to find MAXIMUM energy deposited

- $M_x = 100 \text{ GeV} = 1.8 \times 10^{-25} \text{ kg}$

- $v = 220 \text{ km/s} = 2.2 \times 10^5 \text{ m/s}$

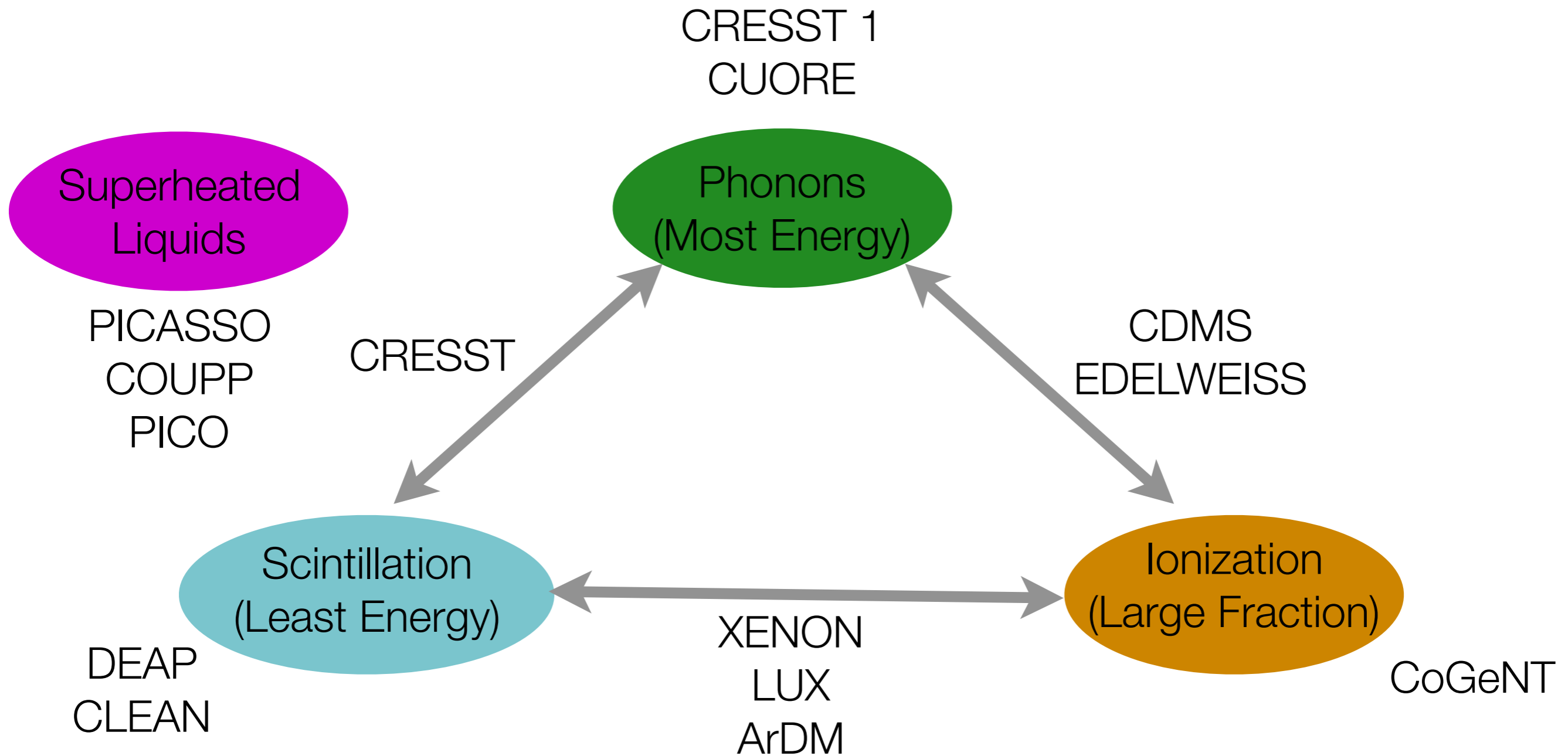
$$T = \frac{1}{2}mv^2 = \frac{1}{2} * 1.8 \times 10^{-25} * (2.2 \times 10^5)^2 = 4.4 \times 10^{-15} \text{ J}$$

- Equivalent to a mosquito flying at 0.00015 kph

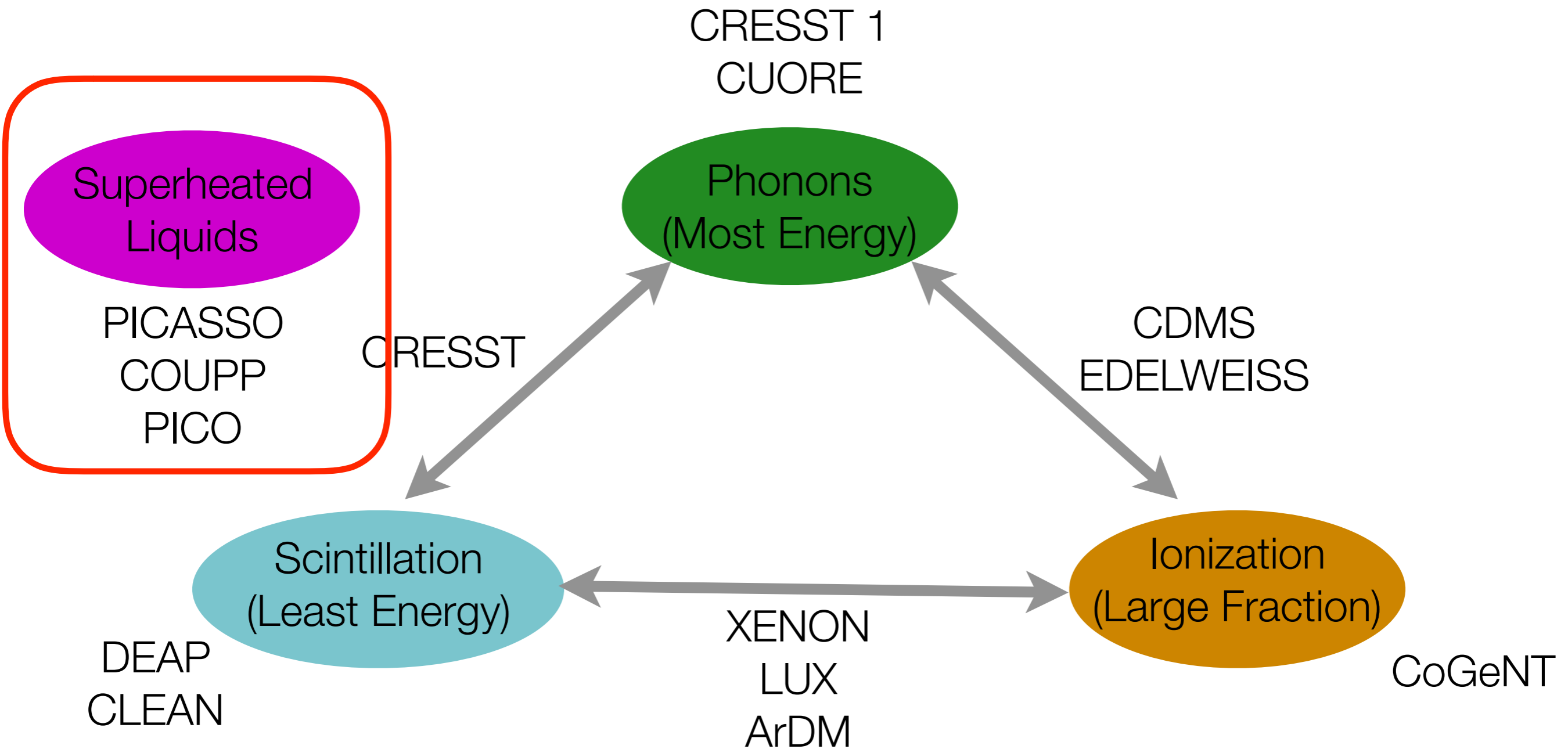
# Background Removal

- Because of this small energy deposit, DM experiments are all about dealing with backgrounds
- Done in two ways
  1. Don't have interactions you don't want
  2. Tell the “bad” interactions from the “good”

# Background Removal

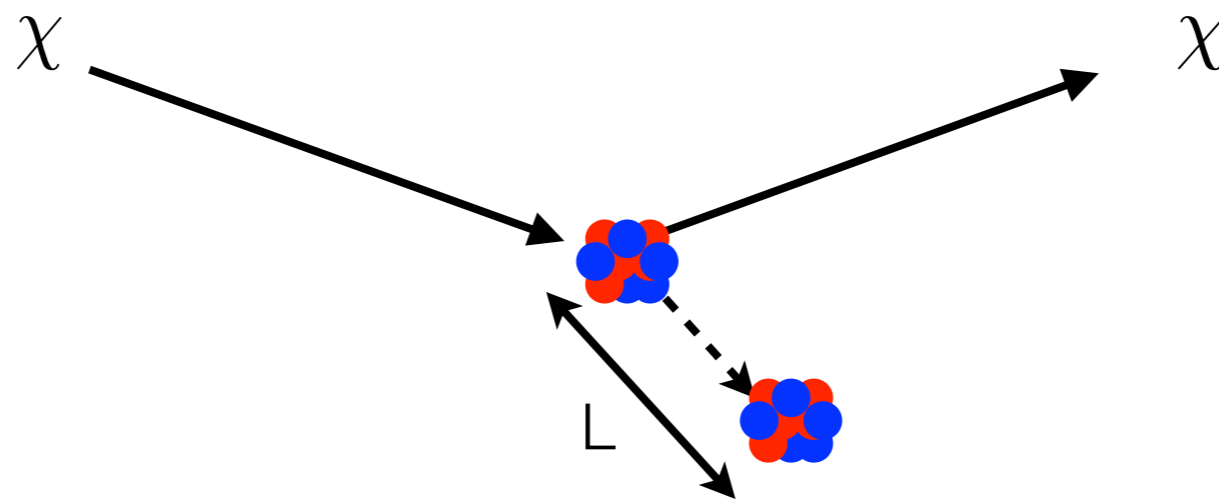


# Background Removal



# PICASSO + COUPP = PICO

- PICO is a novel detector using superheated liquid to amplify the energy deposit
- Small deposit of energy triggers the formation of large bubbles, detectable using acoustic or visual methods





# Seitz Model

- The currently used model says that the energy for the formation of the bubble must come from the interaction, not the surrounding fluid
- This requires a threshold energy deposit in a critical radius

$$E_{threshold} = 4\pi r_c^2 \left( \sigma - T \frac{d\sigma}{dT} \right) + \frac{4\pi}{3} r_c^3 \rho_b (h_b - h_l) - \frac{4\pi}{3} r_c^3 (P_b - P_l)$$

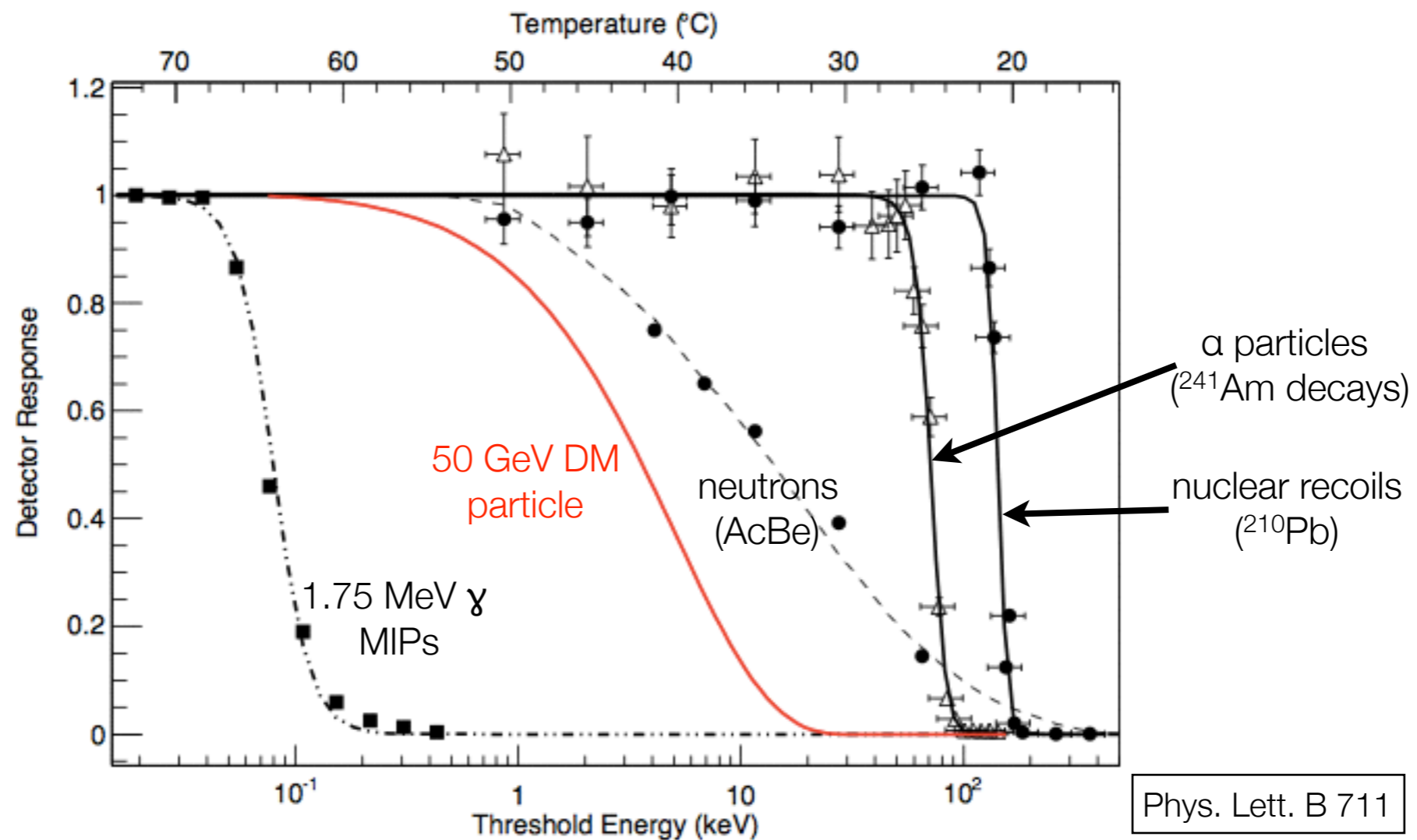
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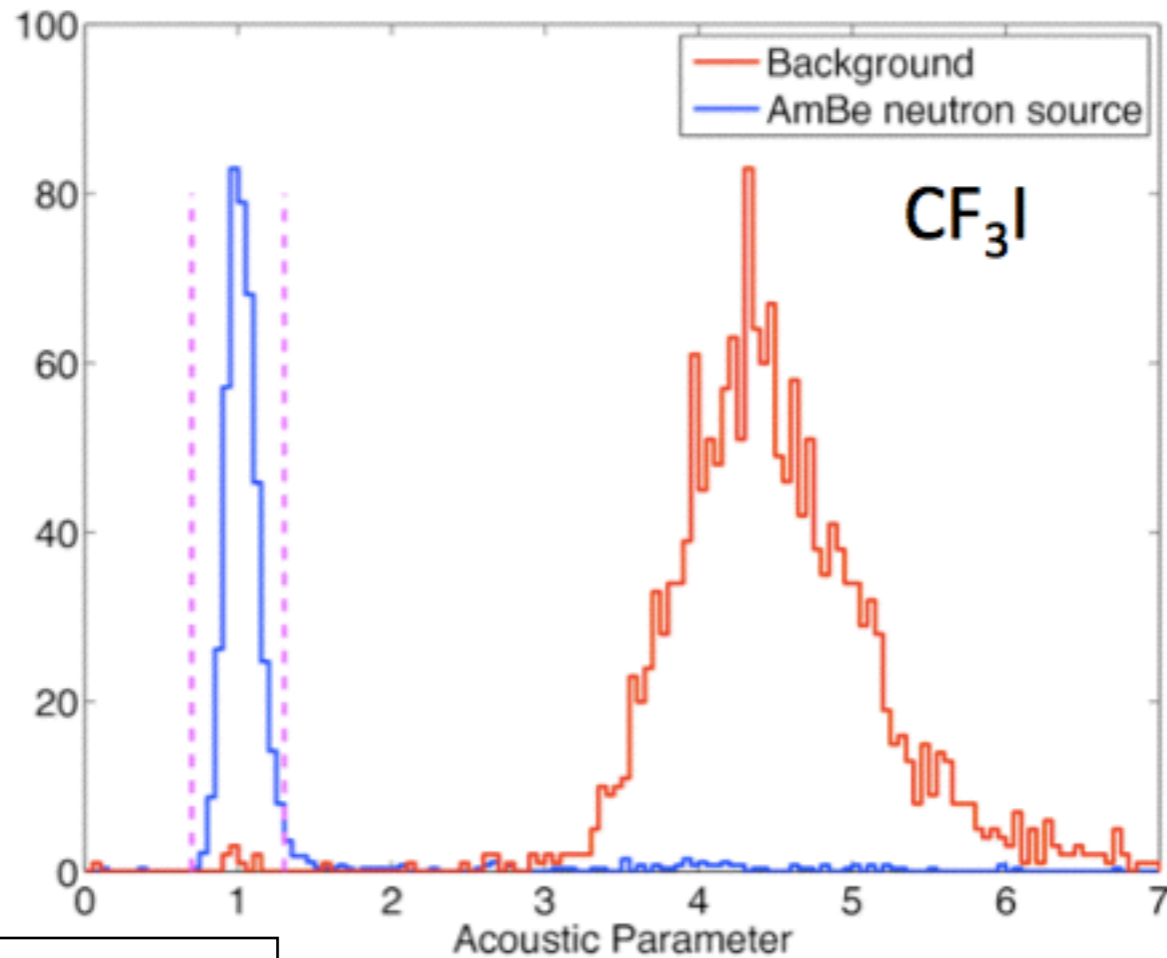
$$E_{threshold} = \underbrace{4\pi r_c^2 \left( \sigma - T \frac{d\sigma}{dT} \right)}_{\text{Overcoming surface tension}} + \underbrace{\frac{4\pi}{3} r_c^3 \rho_b (h_b - h_l)}_{\text{Vaporization of fluid}} - \underbrace{\frac{4\pi}{3} r_c^3 (P_b - P_l)}_{\text{Double counting}}$$

# Backgrounds Method 1

- Gammas and betas are effectively not detected by the detector as they do not meet the  $E_{\text{threshold}}$  in  $r_c$  requirement.



# Backgrounds Method 2



PRD 86:052001

- Distinguish the bubbles formed by backgrounds from those formed by recoil events
- In this case, alphas vs neutrons

$$AP = A(T) \sum_j G_j \sum_n C_n(\vec{x}) \sum_{f_{min}^n}^{f_{max}^n} f \times psd_f^j$$

Temperature Correction  $\rightarrow$   $A(T)$    
 Transducer gain  $\rightarrow$   $G_j$    
 Position correction  $\rightarrow$   $C_n(\vec{x})$    
 Frequency  $\rightarrow$   $f$    
 Power Spectral Density  $\rightarrow$   $psd_f^j$

# The Ultimate Goal (?)

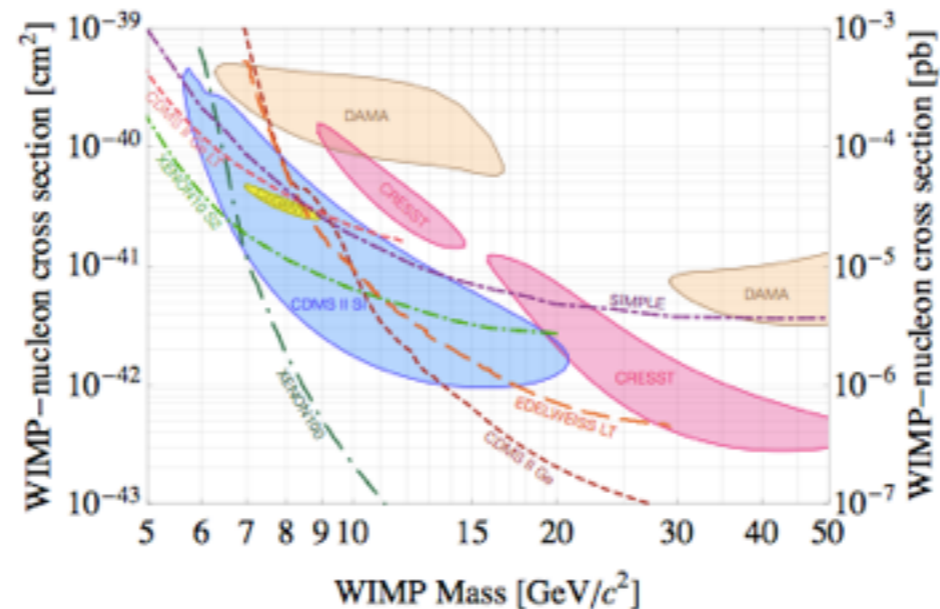


Figure 8. Expanded plot showing spin-independent WIMP-nucleon cross section limits, including closed contours showing hints for low-mass WIMP signals.

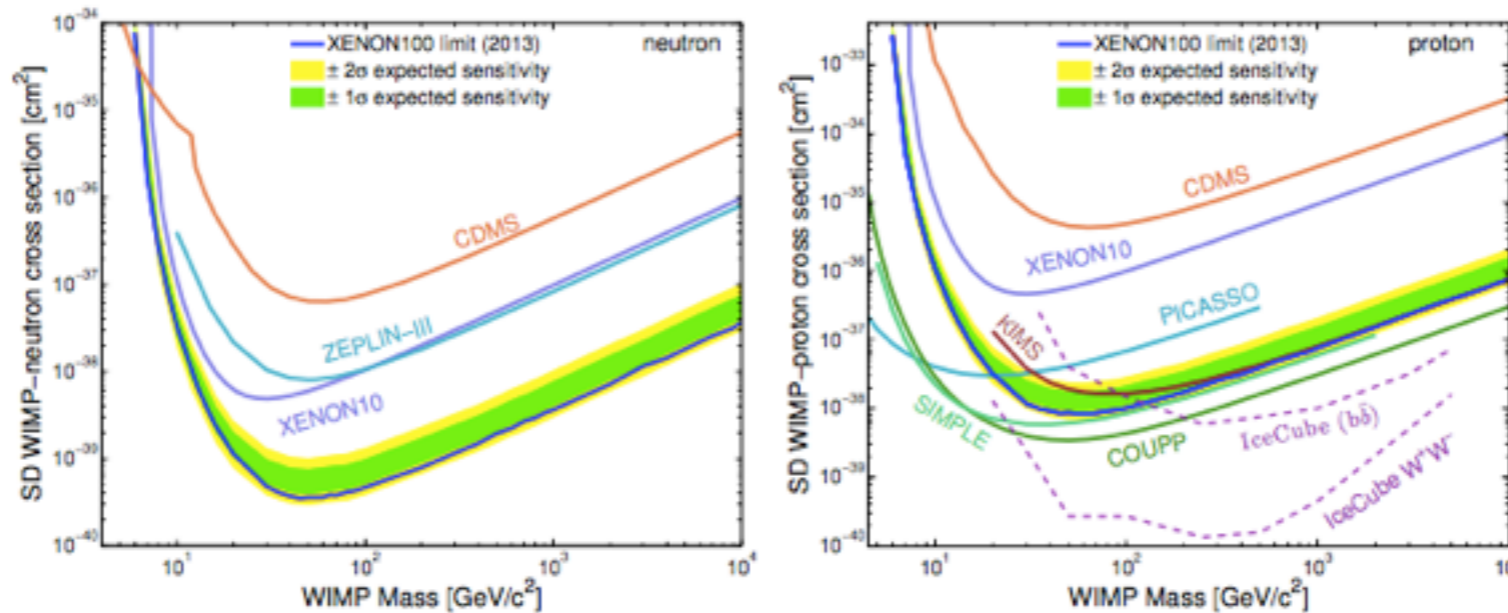


Figure 9. Spin-dependent WIMP-neutron (left) and WIMP-proton (right) cross section limits versus WIMP mass for direct detection experiments [27, 28, 33, 38, 39, 40, 41], compared with the model-dependent Ice Cube results (model-dependent) as of summer 2013 [42].

Snowmass 2013

- Goal should be discovery...

# Calculation of Limit

- Calculation based on 4 parameters
  - Background rate  $B$
  - Background misidentification  $\beta$
  - Signal acceptance  $\alpha$
  - exposure  $MT$  (mass x time)

# Calculation of Limit

- With NO discrimination ( $\beta=1$ ) all interactions are potentially DM
- If there are no events, the 90% CL

$$\begin{aligned}P_{n(obs)}(n_s, 90) &= \frac{(n_s, 90)^{n_{obs}}}{n_{obs}!} e^{-n_s, 90} \\ &= \frac{(n_s, 90)^0}{0!} e^{-n_s, 90} = 0.1 \\ &\rightarrow n_{s, 90} = -\ln(0.1) = 2.3\end{aligned}$$

# Calculation of Limit

- With NO discrimination ( $\beta=1$ ) all interactions are potentially DM
- If there are no events, the 90% CL

$$S_{90} = \frac{2.3}{\alpha MT}$$

- Obviously scales with the exposure time



# Calculation of Limit

- With far more events than expected signal, assume all events are background

$$S_{90} = \frac{N_{BG} + 1.28\sqrt{N_{BG}}}{\alpha MT}$$

- This can be expressed as

$$S_{90} = \frac{\beta}{\alpha} + \frac{1.28}{\alpha} \sqrt{\frac{\beta B}{MT}}$$

# Calculation of Limit

Zero Events

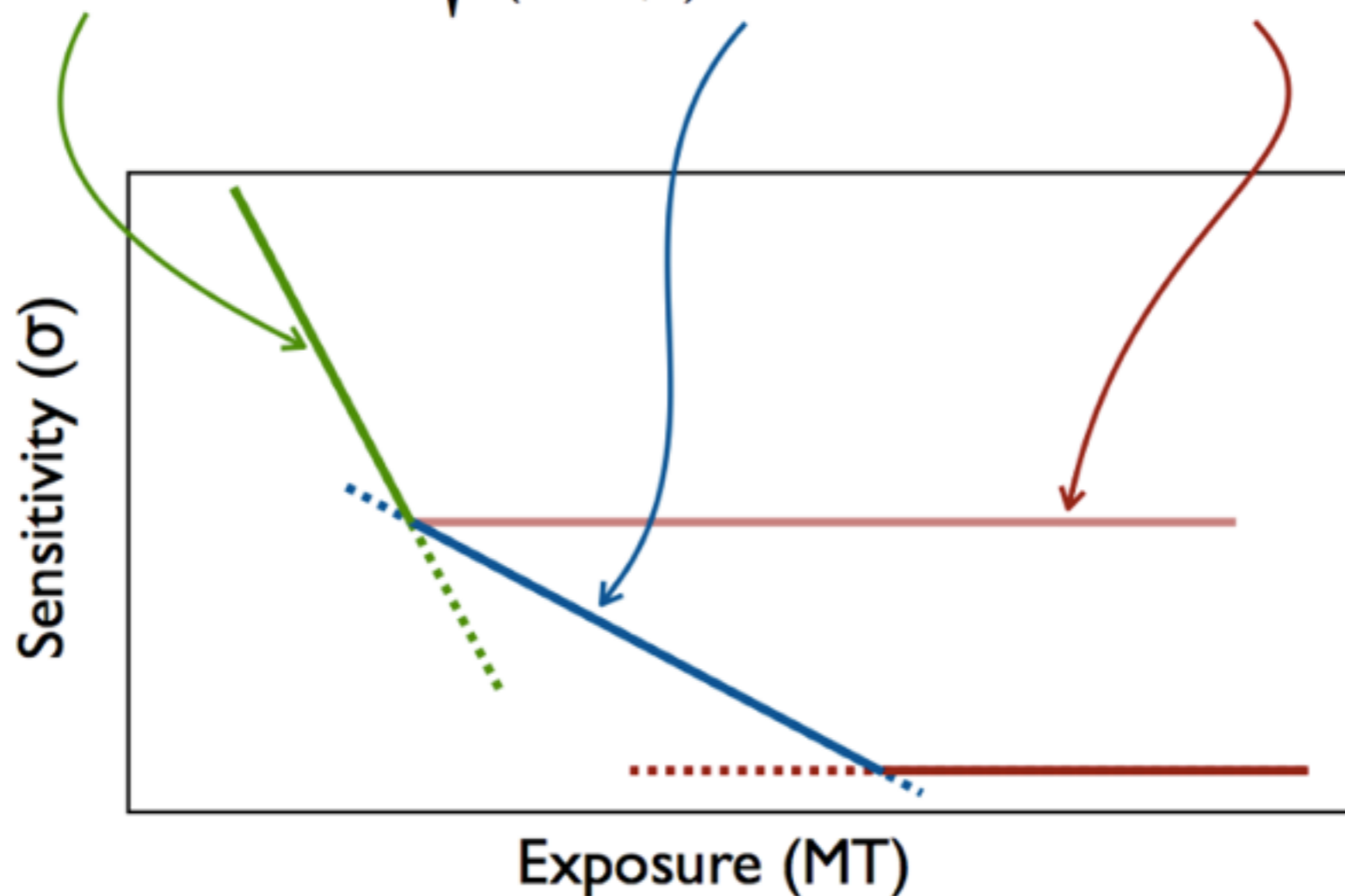
$$S_{90} \propto \frac{2.3}{\alpha MT}$$

Discriminating Detector

$$\sqrt{\frac{\beta(1-\beta)}{(\alpha-\beta)^2}} \sqrt{\frac{B}{MT}}$$

Non-Discriminating

$$\frac{B}{\alpha} + \frac{1.28}{\alpha} \sqrt{\frac{\beta B}{MT}}$$



# A Better Way?

- Let's study a bit more complex example

CRESST 1  
CUORE

Superheated  
Liquids

Phonons  
(Most Energy)

PICASSO  
COUPP  
PICO

CRESST

CDMS  
EDELWEISS

Scintillation  
(Least Energy)

Ionization  
(Large Fraction)

XENON  
LUX  
ArDM

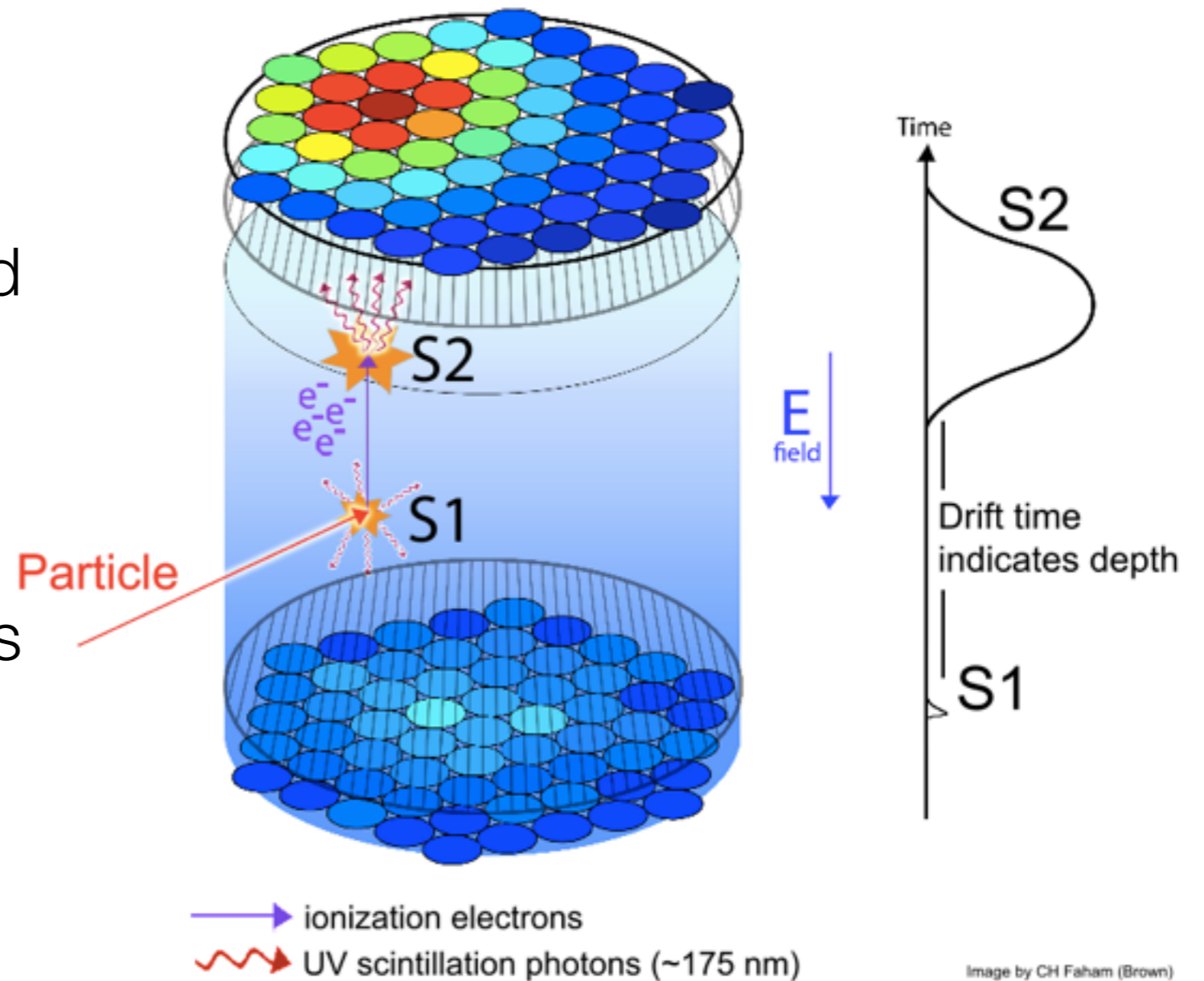
DEAP  
CLEAN

CoGeNT



# Scintillation and Ionization

- Example used here is LUX
- Xenon used in both liquid and gas state within an electric field to amplify small deposit
- Two different scintillations detected, the ratio of which discriminates signal from background



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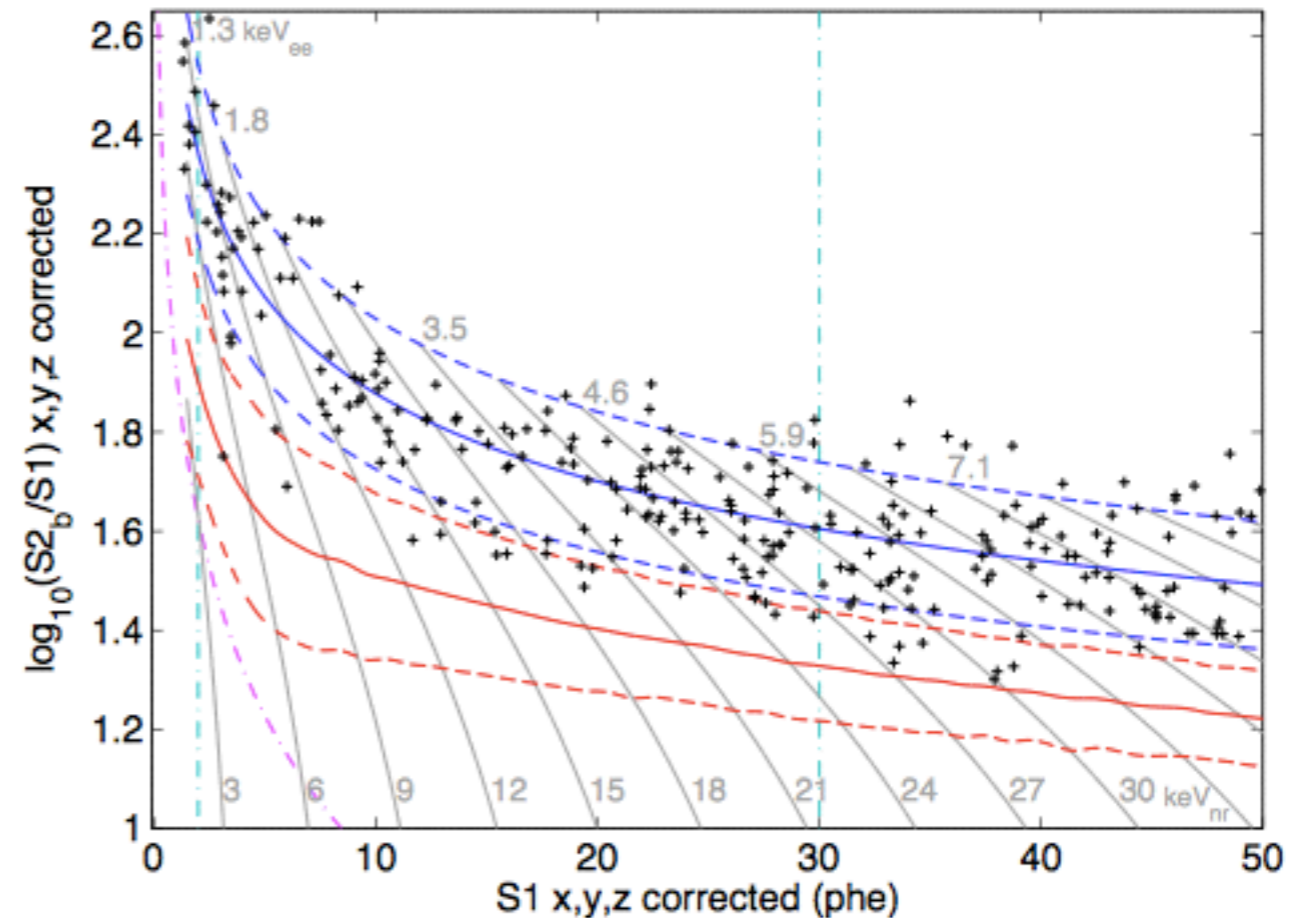
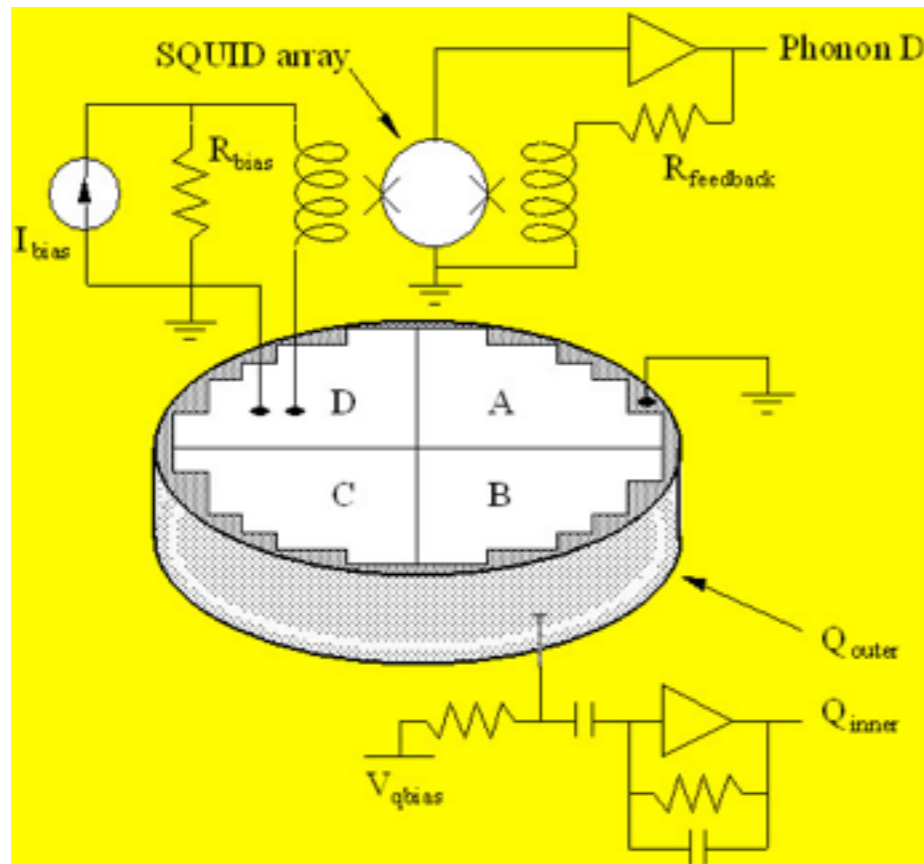


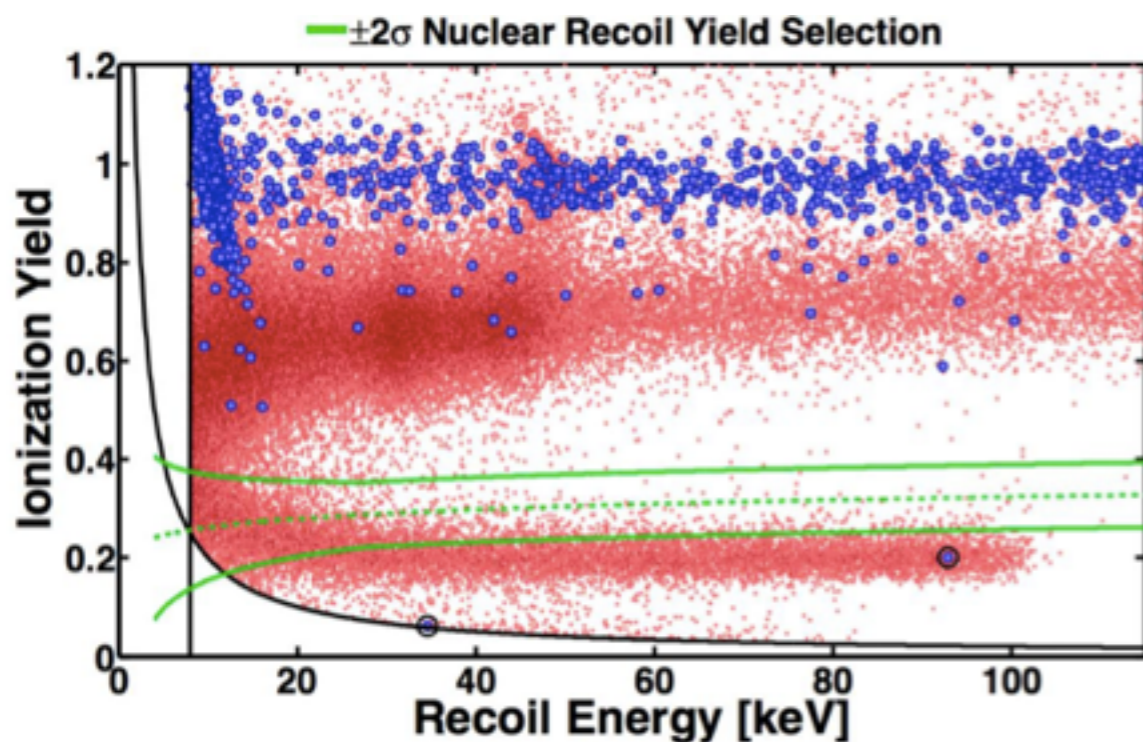
FIG. 4. The LUX WIMP signal region. Events in the 118 kg fiducial volume during the 85.3 live-day exposure are shown. Lines as shown in Fig. [3], with vertical dashed cyan lines showing the 2-30 phe range used for the signal estimation analysis.

arXiv:1310.8214

# Phonons and Ionization



- CDMS similarly uses two channels to distinguish background from signal
- Collecting both ionization and phonons allows for discrimination



# Risks

- The main (and obvious) risk to removing backgrounds is to the lifetime of the experiment
- Aggressive cleaning puts you back into the region discussed previously

# Setting a Limit

- Need to know the number of counts and the distribution of the probability function
- Also need to know the expected number of counts seen



# Expected Counts

- The simplest expectation is:

$$\frac{dR}{dE_R} = \frac{R_0}{E_0 r} e^{-E_R/E_0 r}$$

- $E_R$  is the recoil energy
- $R_0$  is the total event rate
- $E_0$  is the most probable dark matter energy

$$r = \frac{4M_D M_T}{(M_D + M_T)^2}$$

# Trick is in the $R_0$

- Define the  $R_0$  so that it can be calculated
- It's the event rate per unit mass for the earth velocity is 0 and the escape velocity is infinite

$$R_0 = \frac{2}{\pi^{1/2}} \frac{N_0}{A} \frac{\rho_D}{M_D} \sigma_0 v_0$$

# Defining The Bounds

- Really need to define the bounds on the observed number of events
- This has to take into account the backgrounds and the uncertainty associated with those backgrounds
- Three methods here:
  1. Feldman Cousins
  2. Yellin
  3. Binned Likelihood

# Feldman Cousins

- Want to know the confidence region for the number of signal events given the number of observed and background

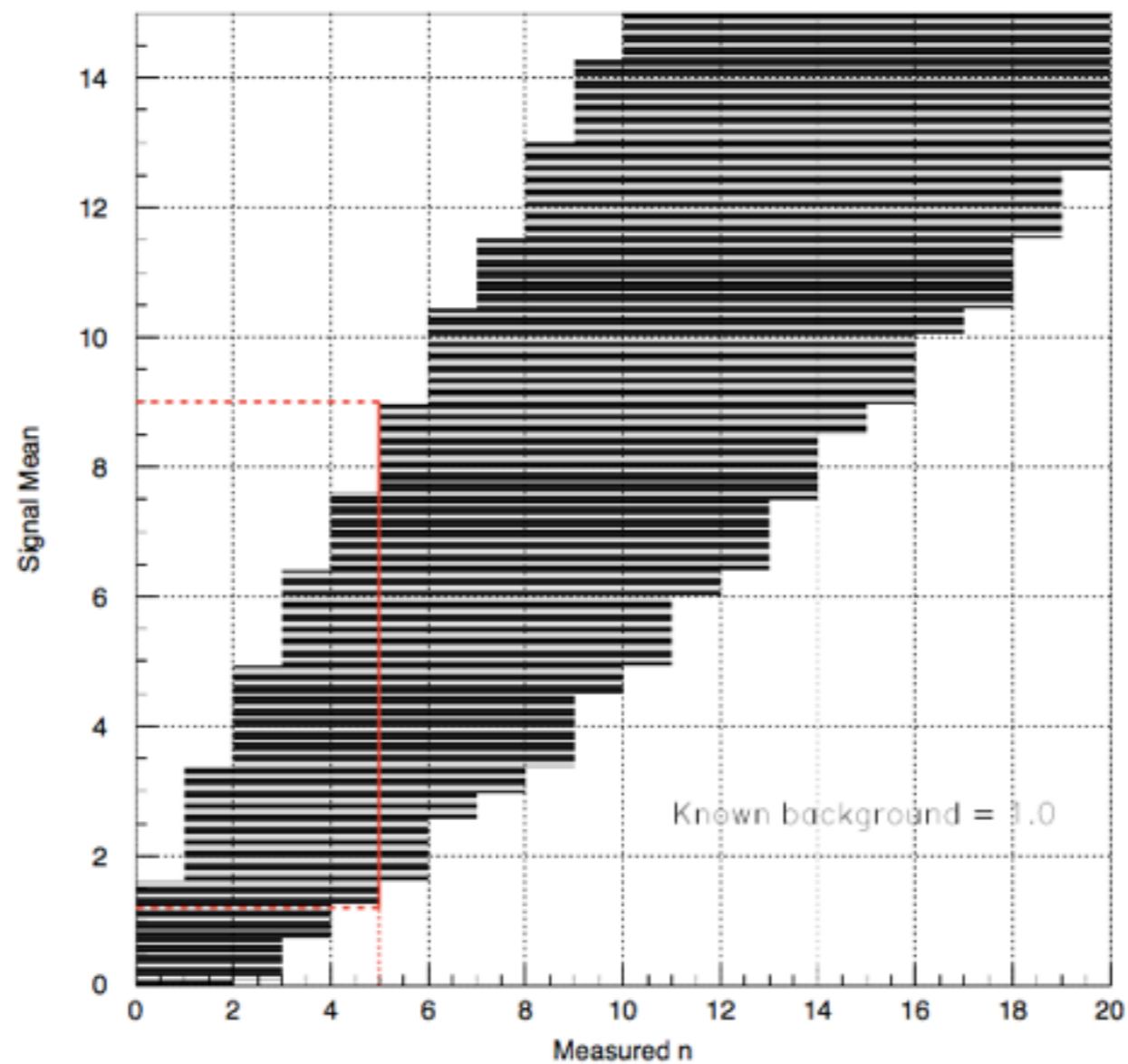
$$P(n|\mu) = \frac{(\mu + b)^n e^{-(\mu+b)}}{n!}$$

- Then maximize the probability, changing  $\mu_{\text{best}}$  and the ratio is the parameter used to define 90%

$$R = \frac{P(n|\mu)}{P(n|\mu_{\text{best}})}$$

# Feldman Cousins

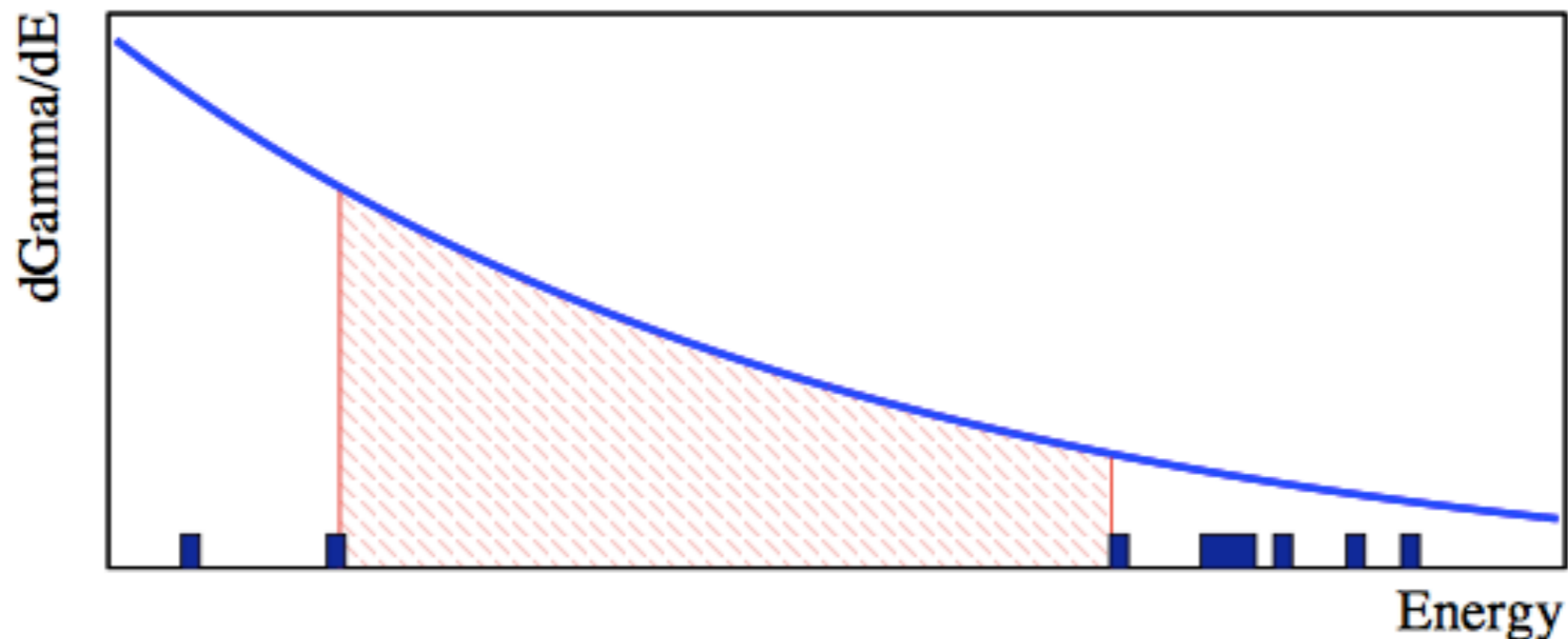
- Plot this and read the results for your experiment!



# But... I'm not sure about my Backgrounds

- All DM experiments are in a new region of detector physics
- The backgrounds are not completely understood
- The best measurement of the backgrounds is the dark matter data itself...

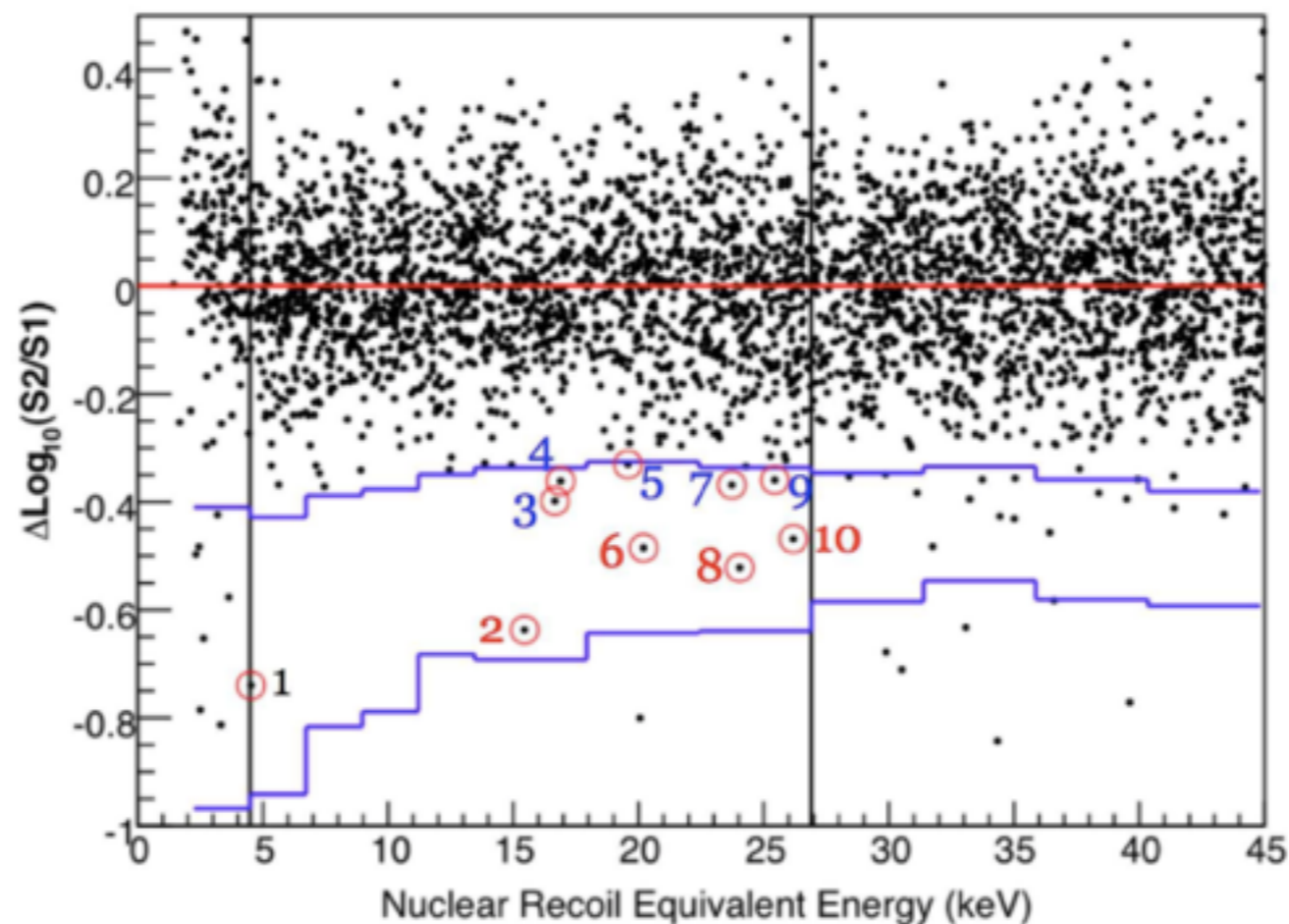
# Yellin to the Rescue!



- The events don't match the expectation well
- Use the “maximum gap” and find the cross-section at which 90% of the trials have a gap smaller than this

# More Yellin

- This can only set an upper limit, and can never be used for discovery
- Also generates one-sided (upper) limits
- No information about the background goes into this (it may be unknown)





# Binned Maximum Likelihood

- We want to use all the information
- Assume we have one discriminator
- Bin the number of counts in that parameter

$$\mathcal{L} = \prod_{i=1}^k P(n_i | \mu(x))$$

$$\ln \mathcal{L} = \sum_{i=1}^k \ln(P(n_i | \mu(x)))$$

# Binned Maximum Likelihood

$$\ln \mathcal{L} = \sum_{i=1}^k \ln(P(n_i | \mu(x)))$$

Use Poisson for low statistics (as before)

$$\ln \mathcal{L} = \sum_{i=1}^k \ln\left(\frac{e^{-\mu_i} \mu_i^{n_i}}{n_i!}\right)$$

But introduce a new term for signal and background counts

$$\mu_x = S \cdot P_s(x) + B \cdot P_b(x)$$

# Binned Maximum Likelihood

$$\mu_x = S \cdot P_s(x) + B \cdot P_b(x)$$

S and B are the hypothetical number of signal and background events in bin x

The likelihood can now be minimized to produce the best estimate of signal and background, which is used to set the  
limit

# Let's Look at Data

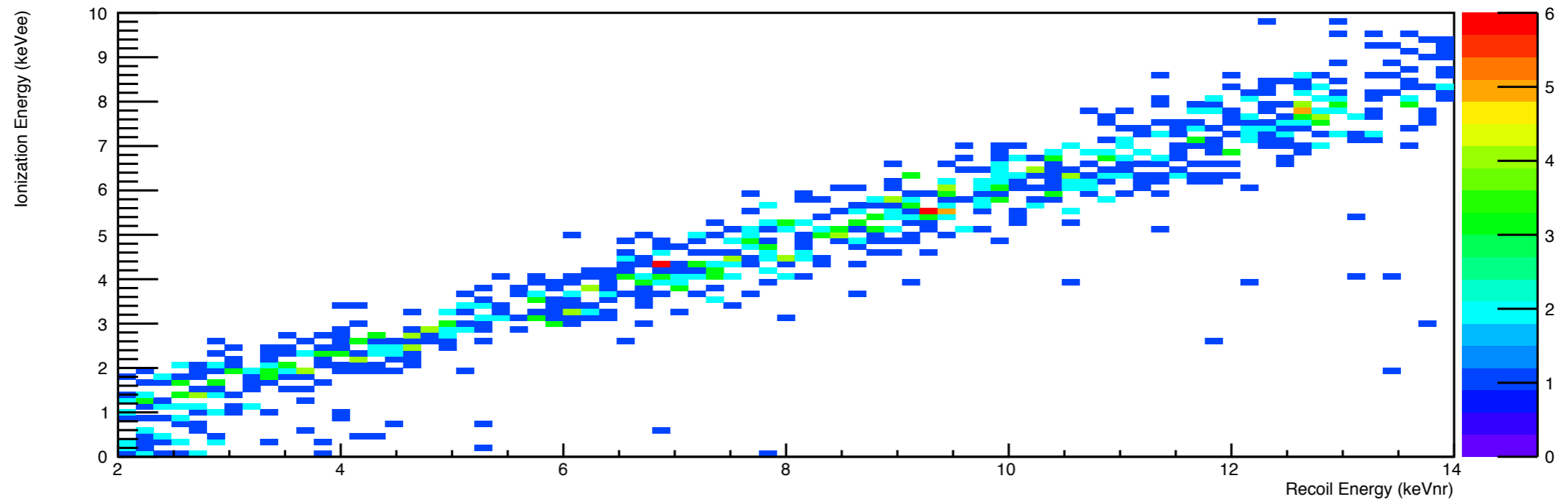
- This is CDMS data taken when I was a postdoc with them
- We have the single and multiple rates for one detector
- Use these to define backgrounds and signal

# Don't have DM Data... as such

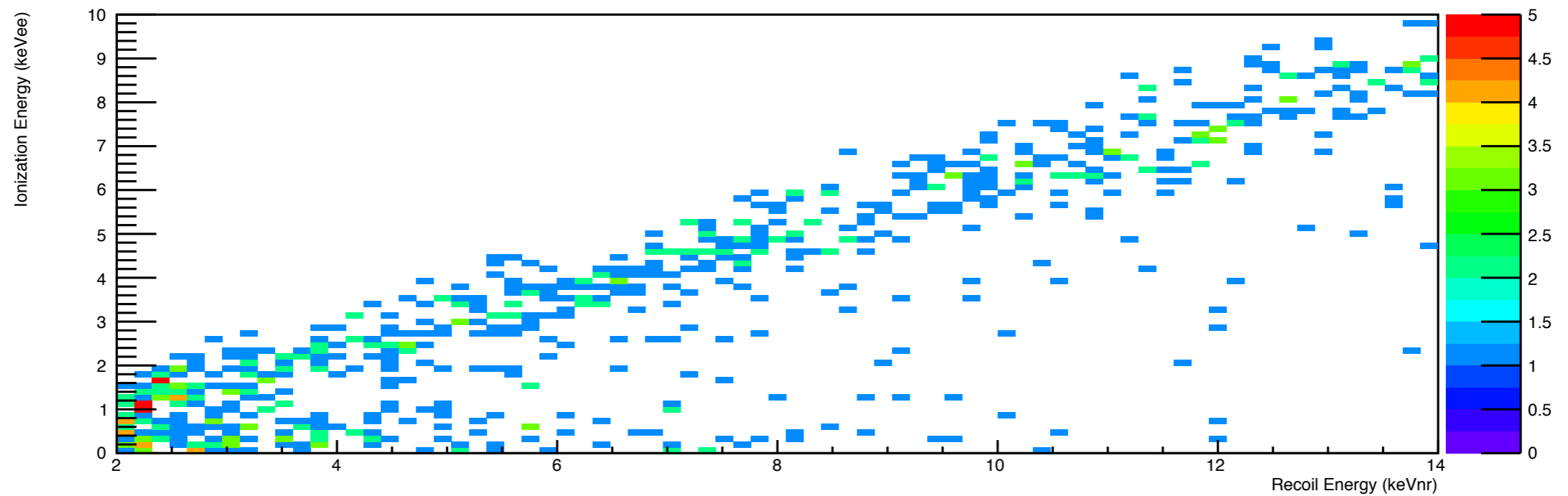
- I have calibration data
- Split into “singles” and “multiples”
- Neutron events will happen in the “singles”
- Let's define “singles” as calibration, “multiples” as data

# CDMS Data

Multiples



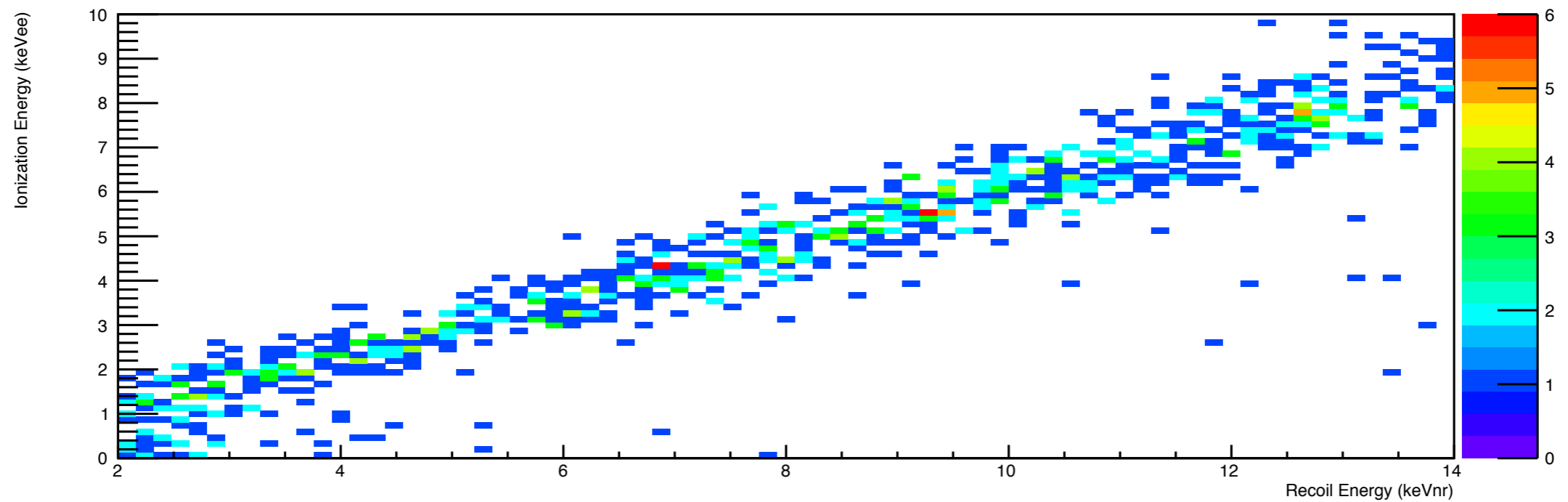
Singles



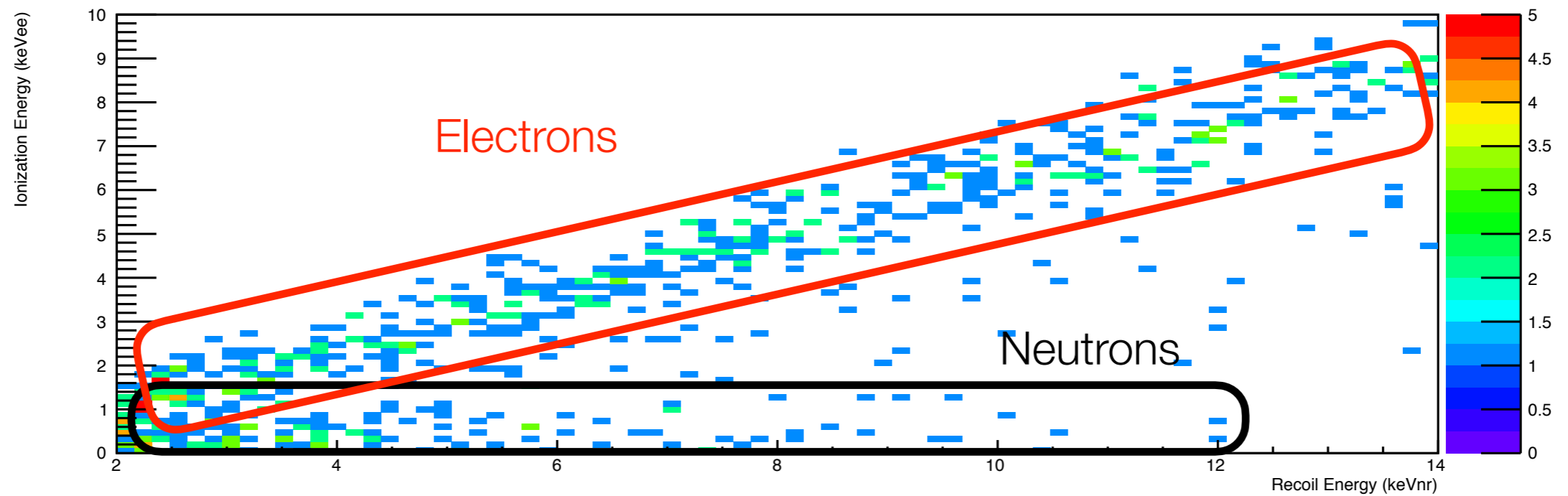
600g Germanium crystal, 1 day exposure

# CDMS Data

Multiples



Singles



600g Germanium crystal, 1 day exposure

Go!