

# On timelike and spacelike minimal surfaces in $AdS_n$ and the Alday-Maldacena conjecture

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- Motivation
- General formalism
- Spacelike minimal surfaces
- No flat spacelike minimal surfaces in  $AdS_n$  beyond 4-cusp case

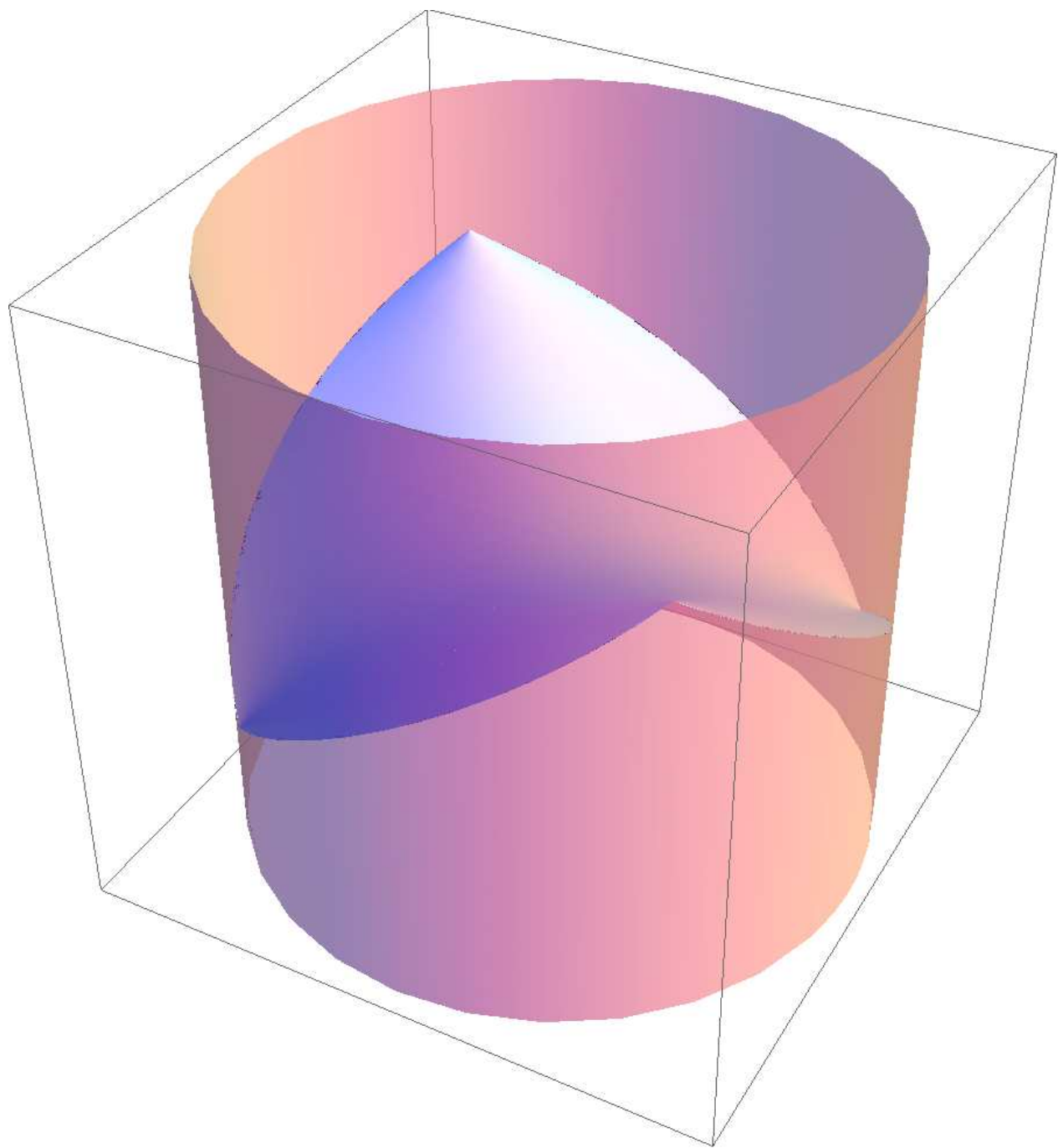
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Alday, Maldacena :

gluon scatt. ampl.  $\Leftrightarrow$  string world surfaces in  $AdS_5$  appr. a lightlike polygon on the boundary

- First lightlike cusp between infinite straight lines in boundary of Poincare patch (= Minkowski<sub>2</sub>)
- extension to full  $AdS_3$  gives totally symmetric tetragon on  $\partial(AdS_3)$
- after isometry trafo in  $AdS_5$  one can reach each lightlike tetragon in Minkowski<sub>4</sub>

But needed minimal surfaces known only for tetragon,  
difficult Plateau like problem



$Y^N$  coordinates in  $\mathbb{R}^{2,n-1}$ ,

$AdS_n$  embedded as hyperboloid

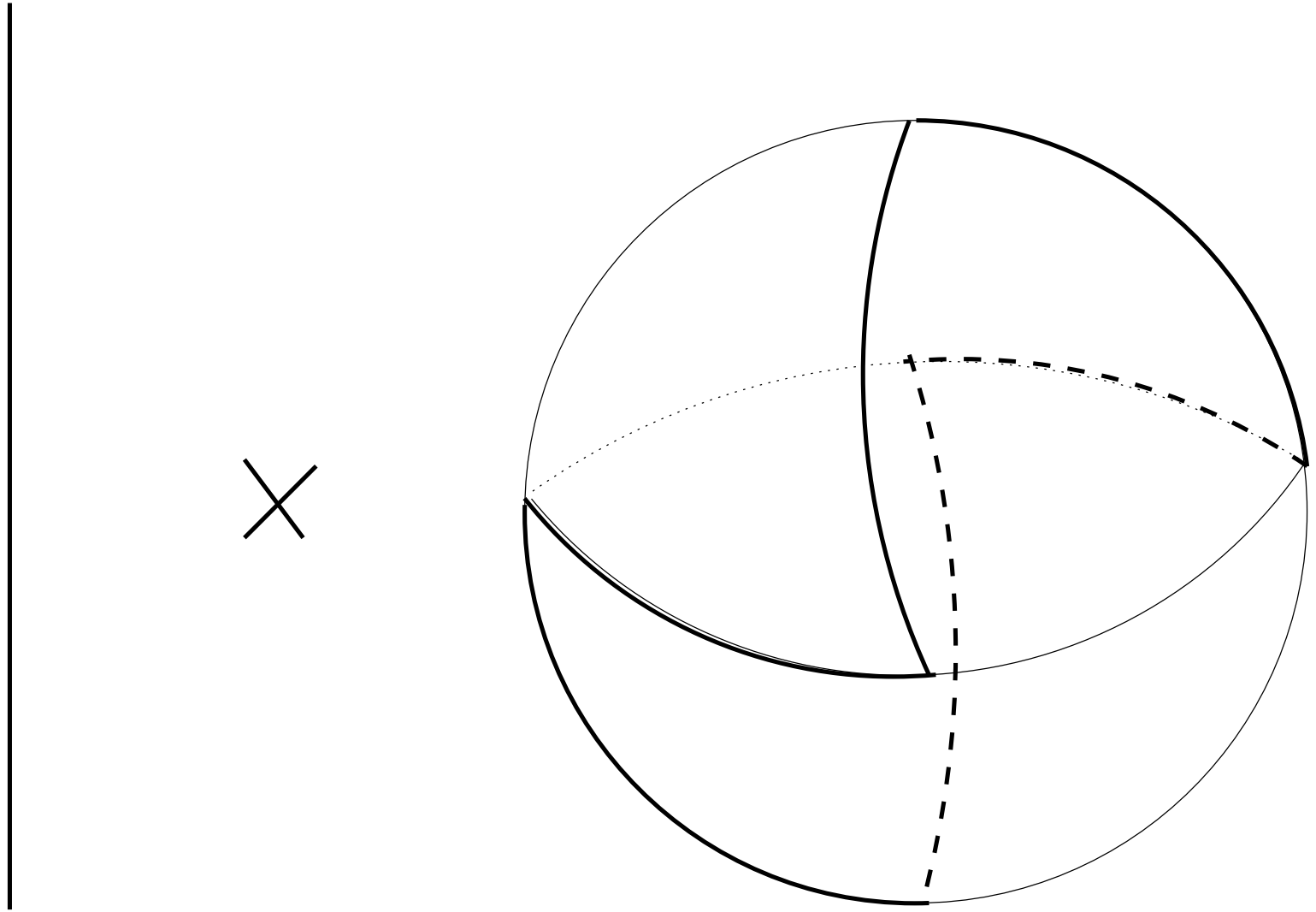
$$(Y^0(X))^2 + (Y^{0'})^2 - (Y^1)^2 - \dots - (Y^{n-1})^2 = 1$$

Tetragon surface in  $AdS_3$ :

$$\begin{aligned} Y^0 &= \cosh \sigma \cosh \tau, & Y^{0'} &= \sinh \sigma \sinh \tau, \\ Y^1 &= \sinh \sigma \cosh \tau, & Y^2 &= \cosh \sigma \sinh \tau. \end{aligned}$$

- induced metric:  $g_{\mu\nu} = \delta_{\mu\nu} \Rightarrow$  minimal, spacelike, flat
- Is there also such surface in max. symmetric situation for  $AdS_4$  or  $AdS_5$ ?

Most symmetric polygon for AdS\_4



spacelike tetragon surface is double Wick rotation of rigid rotating string

$$\tau \mapsto i\tau, \quad Y^2 \mapsto iY^{0'}, \quad Y^{0'} \mapsto iY^2$$

$\exists$  flat timelike minimal surface in  $AdS_5$ :

rigid string with 2 independent rotations (Frolov, Tseytlin)

$$(A \cos \kappa\tau, A \sin \kappa\tau, B \sin \sigma \cos \omega\tau, B \sin \sigma \sin \omega\tau, B \cos \sigma \cos \omega\tau, B \cos \sigma \sin \omega\tau)$$

with  $g_{\mu\nu} = B^2 \eta_{\mu\nu}$

$$\& \quad A^2 - B^2 = 1, \quad \omega^2 = 1 + \kappa^2, \quad B^2 = \frac{\kappa^2}{2}$$

But now to compensate  $\tau \mapsto i\tau$ , one has to Wick rotate  $Y^{0'}, Y^2, Y^4$

$\Rightarrow$  ends up not in  $AdS_5$ .

min. surface (all mean curvatures = 0, stationarity of volume functl.)

$$g^{\mu\nu} \left( \nabla_\mu \partial_\nu X^k(z) + \partial_\mu X^j \partial_\nu X^l \Gamma_{jl}^k(X(z)) \right) = 0 ,$$

$AdS_n$  as a hyperboloid in  $\mathbb{R}^{2,n-1}$ , conformal coord. on the surface

$$\partial \bar{\partial} Y^N(X(z)) - \partial Y^K \bar{\partial} Y_K Y^N = 0 .$$

- $\partial = \partial_\sigma + \partial_\tau, \quad \bar{\partial} = \partial_\sigma - \partial_\tau$  timelike surfaces
- $\partial = \partial_\sigma - i\partial_\tau, \quad \bar{\partial} = \partial_\sigma + i\partial_\tau$  spacelike surfaces

complete the vectors  $Y, \partial Y, \bar{\partial} Y$  to a basis in  $\mathbb{R}^{2,n-1}$  (Pohlmeyer, de Vega/Sanchez, Jevicki et al ....)

$$\{e_N\} = \{Y, \partial Y, \bar{\partial} Y, B_4, \dots, B_{n+1}\} .$$

timelike surfaces  $\Rightarrow$  normal space Euclidean

spacelike surfaces  $\Rightarrow$  normal space Lorentzian

Move the basis along the surface

$$\partial e_N = A_N^K e_K, \quad \bar{\partial} e_N = \bar{A}_N^K e_K. \quad (*)$$

- find a suitable parametrization of the dyn. (geom.) degrees of freedom in the matrices  $A$  and  $\bar{A}$
- derive diff. eqs. for the corr. functions from the eq. of motion (minimal surface condition) and the integrability condition
- after solving these diff. eqs., the surface has to be reconstructed by integrating (\*)

$$\alpha(\sigma, \tau) = \log(\partial Y, \bar{\partial} Y)$$

$$u_a(\sigma, \tau) = (B_a, \partial \partial Y), \quad \bar{u}_a(\sigma, \tau) = (B_a, \bar{\partial} \bar{\partial} Y)$$



$$\begin{aligned} \partial Y &= \partial Y \\ \partial \partial Y &= \partial \alpha \partial Y + u^b B_b \\ \partial \bar{\partial} Y &= e^\alpha Y \\ \partial B_a &= -e^\alpha u_a \bar{\partial} Y + A_a{}^b B_b, \end{aligned}$$

as well as the eqs. generated by replacements  $\partial \leftrightarrow \bar{\partial}$ ,  $u_a \rightarrow \bar{u}_a$ ,  $A_a{}^b \rightarrow \bar{A}_a{}^b$ .

integrability  $\Rightarrow$

$$\begin{aligned} \partial \bar{\partial} \alpha - e^{-\alpha} u^b \bar{u}_b - e^\alpha &= 0 \\ \partial \bar{u}_a - A_a{}^b \bar{u}_b &= 0, \quad \bar{\partial} u_a - \bar{A}_a{}^b u_b = 0, \\ e^\alpha (\bar{u}_a u^b - u_a \bar{u}^b) &= \partial \bar{A}_a{}^b - \bar{\partial} A_a{}^b + \bar{A}_a{}^c A_c{}^b - A_a{}^c \bar{A}_c{}^b. \end{aligned}$$

$a, b, c$  run from  $4, \dots, (n+1)$

Gauss, Codazzi-Mainardi, Ricci

- scalar curvature of surface:  $R = -2 e^{-\alpha} \partial \bar{\partial} \alpha$
- Gauss eq.  $M \subset N$ :  $R_{\alpha\beta\mu\nu}^M - R_{\alpha\beta\mu\nu}^N = (l_{\alpha\mu}^a l_{\beta\nu}^b - l_{\alpha\nu}^a l_{\beta\mu}^b) h_{ab}$
- second fund. forms  $l_{\mu\nu}^a = h^{ab} (B_b, \nabla_\mu \partial_\nu X) = h^{ab} (B_b, \partial_\mu \partial_\nu Y)$

minimal:  $g^{\mu\nu} l_{\mu\nu}^c = 0, \quad \forall c$

timelike minimal  $u = a + b, \quad \bar{u} = a - b : \quad l^c = \frac{1}{2} \begin{pmatrix} a^c & b^c \\ b^c & a^c \end{pmatrix}$

spacelike minimal  $u = a + ib, \quad \bar{u} = a - ib : \quad l^c = \frac{1}{2} \begin{pmatrix} a^c & -b^c \\ -b^c & -a^c \end{pmatrix}$

always either  $g_{\mu\nu}$  or  $h_{ab}$  indefinite

The bar now implies complex conjugation !

By a conformal (holomorphic) trafo  $z \mapsto \zeta(z)$ ,  $\bar{z} \mapsto \overline{\zeta(z)}$  :  $u^c u_c = 1$

with  $u_c = a_c + ib_c \Rightarrow a^c a_c - b^c b_c = 1$  ,  $a^c b_c = 0$

After suitable gauge trafo:

- spacelike I  $(b^c b_c > 0)$  ,  $u^c = (0, i \sinh \beta/2, \cosh \beta/2)$
- spacelike II  $(-1 \leq b^c b_c < 0)$  ,  $u^c = (i \sin \beta/2, \cos \beta/2, 0)$
- spacelike III  $(b^c b_c = 0)$  ,  $u^c = (1 + i\beta, 1 + i\beta, 1)$  .

covariant description of subdivision I - III:

$$T = \frac{1}{8 |\det g|} \epsilon^{\alpha\beta} \epsilon^{\mu\nu} \text{tr}(F_{\alpha\beta} F_{\mu\nu})$$

in conformal coordinates:

$$T = \frac{1}{2} e^{-2\alpha} \text{tr} F^2 = e^{-4\alpha} \left( (\bar{u}_a u^a)^2 - (\bar{u}_a \bar{u}^a)(u_b u^b) \right)$$

with Gauss and  $C = (\bar{u}_a \bar{u}^a)(u_b u^b)$

$$R + 2 \pm 2 e^{-2\alpha} \sqrt{C + e^{4\alpha} T} = 0$$

Exceptional cases:  $C = 0$

Non-exceptional cases: after completely fixing coordinates  $C = 1$

$$e^{-4\alpha} = \frac{(R + 2)^2}{4} - T, \quad \alpha \text{ expressed by invariants !!}$$

$\Rightarrow$  for all minimal surfaces in  $AdS_n$ ,  $n \geq 4$

$$\frac{(R+2)^2}{4} - T \geq 0, \quad \text{saturation by the except. cases}$$

timelike case:  $T \leq 0 \Rightarrow$  no subdivision

spacelike case:

- case I :  $0 \leq T < \frac{(R+2)^2}{4},$
- case II :  $T \leq 0,$
- case III :  $T = 0,$  not all  $F_a^b = 0$

case spacelike I

$$A_5^6 = -\frac{i}{2} \partial\beta, \quad A_4^5 = \rho \cosh \frac{\beta}{2}, \quad A_4^6 = i\rho \sinh \frac{\beta}{2}$$

$$\partial\bar{\partial}\alpha - e^{-\alpha} \cosh \beta - e^\alpha = 0,$$

$$\partial\bar{\partial}\beta + (e^{-\alpha} + \rho\bar{\rho}) \sinh \beta = 0,$$

$$(\bar{\rho}\partial\beta - \rho\bar{\partial}\beta) \sinh \frac{\beta}{2} + (\partial\bar{\rho} - \bar{\partial}\rho) \cosh \frac{\beta}{2} = 0,$$

$$(\bar{\rho}\partial\beta + \rho\bar{\partial}\beta) \cosh \frac{\beta}{2} + (\partial\bar{\rho} + \bar{\partial}\rho) \sinh \frac{\beta}{2} = 0.$$

case spacelike II

$$A_4^5 = \frac{i}{2} \partial\beta, \quad A_4^6 = \rho \cos \frac{\beta}{2}, \quad A_5^6 = i\rho \sin \frac{\beta}{2}.$$

$$\partial\bar{\partial}\alpha - e^{-\alpha} \cos \beta - e^\alpha = 0,$$

$$\partial\bar{\partial}\beta + (e^{-\alpha} + \rho\bar{\rho}) \sin \beta = 0,$$

$$(\bar{\rho}\partial\beta - \rho\bar{\partial}\beta) \sin \frac{\beta}{2} - (\partial\bar{\rho} - \bar{\partial}\rho) \cos \frac{\beta}{2} = 0,$$

$$(\bar{\rho}\partial\beta + \rho\bar{\partial}\beta) \cos \frac{\beta}{2} + (\partial\bar{\rho} + \bar{\partial}\rho) \sin \frac{\beta}{2} = 0.$$

case spacelike III

$$\partial\bar{\partial}\alpha - 2 \cosh \alpha = 0,$$

.....  
.....

On flat surface always  $\exists$  coordinates with metric  $\eta_{\mu\nu}$  or  $\delta_{\mu\nu}$ .

Here the diffeom. freedom fixed already,  $\Rightarrow$  to analyze all  $\partial\bar{\partial}\alpha = 0$

$AdS_3$ : From sinh-Gordon  $\Rightarrow \alpha = 0$

$$A_N^K = \begin{pmatrix} 0 & 1 & 0 & 0 \\ 0 & 0 & 0 & -i \\ 1 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 \end{pmatrix}, \quad \bar{A}_N^K = \begin{pmatrix} 0 & 0 & 1 & 0 \\ 1 & 0 & 0 & 0 \\ 0 & 0 & 0 & i \\ 0 & i & 0 & 0 \end{pmatrix}.$$

Reconstruction of surface: Solve linear first order eqs. for frame  $\{e_N\}$

$$e_N(\sigma, \tau) = \mathcal{M}_N^K e_K(0, 0), \quad \mathcal{M}_N^K = \left( \exp\left(\frac{\sigma + i\tau}{2} A\right) \exp\left(\frac{\sigma - i\tau}{2} \bar{A}\right) \right)_N^K.$$

Brute force exponentiation by Mathematica

Using isomorphy  $SO(2, 2)$  to  $SU(1, 1) \times SU(1, 1)$  more elegant, application to solitonic solutions of sinh-Gordon (i.e. non constant  $\alpha$ ) (Jevicki, Jin, Kalousios, Volovich)



$$\mathcal{M}_N^K = \begin{pmatrix} C_\sigma C_\tau & i \bar{U}_{\sigma,\tau} & -i U_{\sigma,\tau} & S_\sigma S_\tau \\ -i U_{\sigma,\tau} & C_\sigma C_\tau & -i S_\sigma S_\tau & \bar{U}_{\sigma,\tau} \\ i \bar{U}_{\sigma,\tau} & i S_\sigma S_\tau & C_\sigma C_\tau & U_{\sigma,\tau} \\ S_\sigma S_\tau & U_{\sigma,\tau} & \bar{U}_{\sigma,\tau} & C_\sigma C_\tau \end{pmatrix}$$

with

$$C_\sigma = \cosh \frac{\sigma}{\sqrt{2}}, S_\sigma = \sinh \frac{\sigma}{\sqrt{2}}, U_{\sigma,\tau} = \frac{1+i}{2\sqrt{2}} \left( \sinh \frac{\sigma+\tau}{\sqrt{2}} + i \sinh \frac{\sigma-\tau}{\sqrt{2}} \right)$$

string position  $Y(\sigma, \tau)$  is the first vector of the frame  $\{e_N\}$

second and third vector not orthonormalized, take corr. linear comb.

- surface can be read off from first row of matrix  $\mathcal{M}_N^K$
- result is four cusp surface used by Alday-Maldacena
- freedom of overall  $SO(2,2)$  trafo encoded in choice of starting frame  $\{e_N(0,0)\}$

$AdS_5$ : form of  $\alpha$ -equations  $\Rightarrow$  no flat surface in cases spacelike I, III

case spacelike II:  $\partial\bar{\partial}\alpha = 0 \Rightarrow \cos\beta = -e^{2\alpha}$

assume first  $\sin\beta \neq 0 \Rightarrow \partial\bar{\partial}\beta = \frac{4e^{2\alpha}}{\sin\beta} \left(1 - \frac{\cos\beta e^{2\alpha}}{\sin^2\beta}\right) \partial\alpha\bar{\partial}\alpha$

insert in  $\beta$ -eq.  $\Rightarrow 4e^{2\alpha} \left(1 + \frac{e^{4\alpha}}{\sin^2\beta}\right) + (e^{-\alpha} + \rho\bar{\rho}) \sin^2\beta = 0$

$\Rightarrow$  contradiction (note:  $\rho\bar{\rho}$  not pos. def. in timelike case)

$\sin\beta = 0: \Rightarrow$  ( since  $\cos\beta < 0$  )  $\cos\beta = -1$

$\rho, \bar{\rho}$ -equations degenerate to  $\partial\bar{\rho} + \bar{\partial}\rho = 0$

$$A_N^K = \begin{pmatrix} 0 & 1 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & -i & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & -i & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & i\rho \\ 0 & 0 & 0 & 0 & -i\rho & 0 \end{pmatrix}, \quad \bar{A}_N^K = \begin{pmatrix} 0 & 0 & 1 & 0 & 0 & 0 \\ 1 & 0 & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & i & 0 & 0 \\ 0 & i & 0 & 0 & 0 & 0 \\ 0 & 0 & 0 & 0 & 0 & -i\bar{\rho} \\ 0 & 0 & 0 & 0 & i\bar{\rho} & 0 \end{pmatrix}$$

- Both matrices are block diagonal, exponentials, too
- new degrees of freedom relative to the  $AdS_3$  case, encoded in the lower right blocks with  $\rho$  and  $\bar{\rho}$ , have no influence on first row

$\Rightarrow$  All flat spacelike minimal surfaces in  $AdS_5$  are realized in a subspace

$AdS_3 \subset AdS_5$  and are of type just discussed.

Straightforward extension to  $AdS_n$  has been done, too.

- 
- Studied Pohlmeyer type reduction and geometric language for  $AdS_n$
  - Analyzed differences timelike versus spacelike, subclasses of spacelike
  - Exist flat timelike minimal surfaces beyond those in an  $AdS_3$  subspace
  - Exist no flat spacelike minimal surfaces beyond those in  $AdS_3$  subspaces
  - Did also classification of all flat timelike surfaces in  $AdS_5$ .  
All flat timelike minimal surfaces can be parameterized by two arbitrary free chiral functions  $\phi(z)$ ,  $\bar{\phi}(\bar{z})$ ,  $\alpha = \phi + \bar{\phi}$ .
  - Reduction for generic  $n$  with full gauge structure for rotations/Lorentz trafo gives interesting relations to WZW models