

# String Quantization in $d$ dimensions

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## Introduction

Classical Geometry of Strings

Iterated integrals and reparametrization invariance

## Canonical String Quantization

### Algebraic properties of classical strings

The shuffle algebra

Poisson structure

### Invariant Quantization

Quantization: History

Quantization via Kernel of a derivation

**Nambu-Goto action:** Worldsheet  $\Sigma$  is a critical point of  $S = \int d^2 \text{vol} = \text{area of a (2-dim) surface immersed into } (\mathbb{R}^d, \eta)$  with  $\eta$  Euclidean or Minkowskian metric  $\Rightarrow \Sigma$  is a **geometric object**.

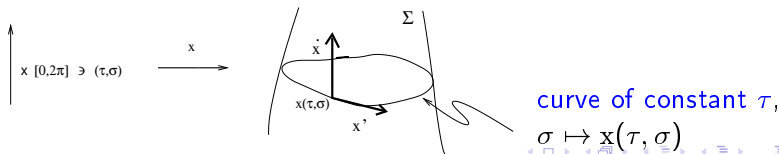
**Parametrization**  $x : \mathbb{R} \times [0, 2\pi] \rightarrow \mathbb{R}^d$   
 $(\tau, \sigma) \mapsto x_\mu(\tau, \sigma), \quad \mu \in \{0, \dots, d-1\}$

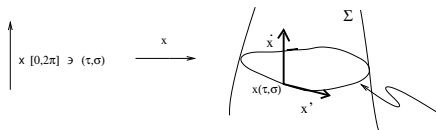
For  $\eta$  Minkowski, choose  $x$  such that  $\dot{x} := \partial_\tau x \in \mathbb{R}^d$  is timelike,  
 $x' := \partial_\sigma x \in \mathbb{R}^d$  is spacelike.

**Lagrangian density** for  $S$  is  $L = \sqrt{-\det g}$ ,  $g$  (induced metric on  $\Sigma$ )

**Momentum**  $p = \frac{\partial L}{\partial \dot{x}}$  In conformal coordinates  $p = \dot{x}$ .

This talk:  $\Sigma$  tube-shaped (closed string).





curve of constant  $\tau$ ,  
 $c : \sigma \mapsto \mathbf{x}(\tau, \sigma)$

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**Hamiltonian Formalism**: Canonical Poisson structure.

**Constraints**  $p^2 + (x')^2 \approx 0$  (lapse),  $p x' \approx 0$  (shift) are **first class**  
 $\rightarrow$  **Gauge symmetry**

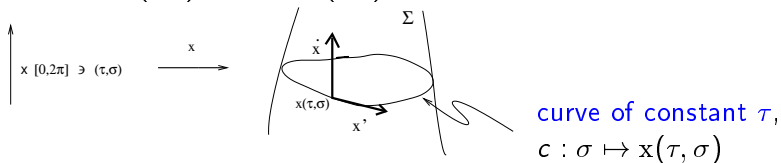
Constraints are infinitesimal generators of **reparametrizations**.

**Reparametrization invariant** objects Poisson-commute with the constraints.

# Reparametrization invariant integrals

**Right/Left mover**  $u^\pm := p \pm x'$ , i.e.  $u^\pm(\tau, \sigma) \in T_{x(\tau, \sigma)} \Sigma \simeq \mathbb{R}^d$ .

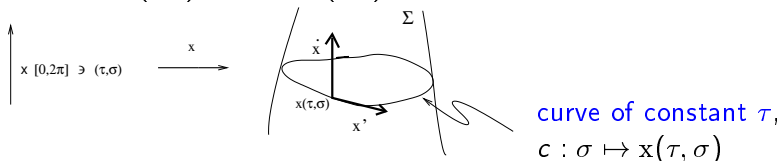
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Define for any  $\sigma \in [0, 2\pi)$ , any  $n \in \mathbb{N}$ , any  $\mu_1, \dots, \mu_n \in \{0, \dots, d-1\}$  the iterated integral along  $c$ , starting in  $c(\sigma)$

$$R_{\mu_1 \dots \mu_n}^\pm(\tau, \sigma) = \int_{\sigma}^{\sigma+2\pi} d\sigma_1 \int_{\sigma}^{\sigma_1} d\sigma_2 \cdots \int_{\sigma}^{\sigma_{n-1}} d\sigma_n u_{\mu_1}^\pm(\tau, \sigma_1) \cdots u_{\mu_n}^\pm(\tau, \sigma_n)$$

where  $u_{\mu}^\pm$ ,  $\mu \in \{0, \dots, d-1\}$  are the  $d$  components of  $u^\pm \in \mathbb{R}^d$ .

[Pohlmeyer, Pohlmeyer+Rehren since 80's]

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Define  $Z_{\mu_1 \dots \mu_n}^\pm := R_{\mu_1 \dots \mu_n}^\pm + R_{\mu_2 \dots \mu_n \mu_1}^\pm + \cdots + R_{\mu_n \mu_1 \dots \mu_{n-1}}^\pm$  (**cyclic sum**)



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**Theorem** Cyclic sums  $Z$  are **independent of the parametrization**  $x$ ,  
**Poisson-commute with the constraints.**

In particular,  $\partial_\sigma Z = \text{const}$   $\partial_\tau Z = 0$ .

**Method:** Nambu-Goto string = **integrable system**.

**Reconstruction** of worldsheet from set of all  $Z$  (complete system).

[Pohlmeyer, Pohlmeyer+Rehren since 80's]

# Canonical framework

**Fourier decomposition** of the string's parametrization  $x$   
 → re-interpretation as **conformal field theory**.

Difficulty: **restore reparametrization invariance** (only consistent in critical dimension, i.e.  $d = 26$ )

In general, a polynomial in Fourier modes  $\alpha$  is not invariant under reparametrizations. But: the invariant cyclic sums  $Z_{\mu_1 \dots \mu_n}$  can be written as linear combinations of infinite sums of polynomials in Fourier modes, e.g. for  $Z_{\mu_1 \dots \mu_4}$

$$\sum_{n,m>0} \frac{1}{nm} \alpha_{\mu_1}^{-n} \alpha_{\mu_2}^{-m} \alpha_{\mu_3}^n \alpha_{\mu_4}^m, \quad \sum_{n,m>0} \frac{1}{nm} \alpha_{\mu_1}^{-n} \alpha_{\mu_2}^{-m} \alpha_{\mu_3}^{n+m} \alpha_{\mu_4}^0 \quad \text{and more}$$

Poisson-commute with constraints, e.g. classical BRST charge...

## BRST approach

BRST Charge  $\Omega_0$ , an infinite sum of 2-fold products of Fourier modes  $\sum \alpha\alpha$  with the property, that for any infinite sum of products of Fourier modes  $X = \sum \alpha \cdots \alpha$ , we have

$$\{\Omega_0, X\} = 0 \quad \text{iff } X \text{ is reparametrization invariant}$$

Observe: identification of reparametrization invariance (= physical relevance) via an **inner derivation!**

**Canonical quantization:** replace positive Fourier modes by annihilation operators, negative Fourier modes by creation operators on Fock space, replace polynomials by normal ordered counterparts (creation operators on the left).

**Problem:** concept not invariant under changes of the parametrization: in general, a reparametrization mixes positive and negative modes.

## BRST approach II

One considers the **BRST operator**  $\Omega$  on Fockspace, a nilpotent operator  $\Omega =: \Omega_0 : + \text{more}$ , with  $: \Omega_0 :$  denoting the canonically quantized counterpart of  $\Omega_0$ .

**Idea:** the physically relevant subspace in Fockspace is  $\ker \Omega$  and physically relevant operators on Fock space commute with  $\Omega$  on  $\ker \Omega$ .

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**Prop** [Bahns] Canonically quantized cyclic sums  $Z_{\dots}$  of the iterated integrals  $R_{\dots}$  do **not** commute with  $\Omega$  on  $\ker \Omega$  **regardless of the dimension**.

**Proof:** obstructions appear for any  $Z_{\mu_1 \dots \mu_n}$  with  $n \geq 3$ . ( $n \leq 2$  is ok)

Observe: in usual approach one considers infinite sums of at most **two-fold** products of operators.

## What does this mean?

Canonical approach corresponds to a **choice of representation**  
(Fock space)

**Alternative:** Invariant Quantization [Pohlmeyer's programme]

Classical **Poisson algebra**, bracket  $\{Z_{\dots}, Z_{\dots}\} = \text{linear comb. of } Z_{\dots}$   
and shuffle multiplication on set of  $Z$ . **Deform** this Poisson algebra  
to an associative **quantum algebra**. As a **last step**, study  
**representation theory** of the quantum algebra.

Attempts at representation/quantization (in one step)

- Fock space: no representation space for algebra of invariants [B].
- Thiemann: lqg framework (no critical dimension)
- Schreiber: DDF operators (problem with Lorentz covariance)

# The shuffle algebra

Results on the invariants  $Z$  due to Pohlmeyer, Rehren and collaborators since 1986. Reformulated within framework of the shuffle algebras over a (finite) alphabet [Bahns 08]  $R_{\mu_1 \dots \mu_n} \rightarrow \text{word } \mu_1 \dots \mu_n$ .

## The shuffle algebra

- ▶ Let  $d \in \mathbb{N}$ , consider the alphabet  $A = \{0, \dots, d - 1\}$  and the set of words  $A^*$ . Let  $W_n$  denote the free module over  $\mathbb{C}$  with basis given by all words  $\in A_n^* =$  the set of words of length  $n$ .



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- ▶ **shuffle product** of words  $x = x_1 \cdots x_n$  and  $y = y_1 \cdots y_m \in A^*$

$$x \# y := x_1(x_2 \cdots x_n \# y) + y_1(x \# y_2 \cdots y_m) \in A_{n+m}^*$$

linear extension  $\Rightarrow (\bigoplus W_n, \#)$  **shuffle algebra**  $Sh(A)$ .

Shuffle product of words  $x = x_1 \cdots x_n$  and  $y = x_{n+1} \cdots x_{n+m} \in A^*$

$$\begin{aligned}\mu(x \otimes y) &:= x \# y = x_1(x_2 \cdots x_n \# y) + x_{n+1}(x \# x_{n+2} \cdots x_{n+m}) \\ &= \sum_{\pi \in S_{n,m}} x_{\pi(1)} \cdots x_{\pi(n+m)}\end{aligned}$$

$S_{n,m} = S_{n+m}/S_n \times S_m =$  set of shuffle permutations  
(mixing two decks of cards).

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**Deconcatenation coproduct**  $\Delta : Sh(A) \rightarrow Sh(A) \otimes Sh(A)$

$$\Delta(x) = \sum_{I \sqcup J = \underline{n}} x_I \otimes x_J \quad \text{for } x \in W_n$$

where  $\underline{n} = \{1, \dots, n\}$ ,  $I \sqcup J = \underline{n}$  as ordered sets (including  $I, J = \emptyset$ ), and for  $I = \{1, \dots, k\}$  we set  $x_I = x_1 \cdots x_k$ , and  $x_\emptyset = 1$ . Then  $(Sh(A), \mu, \Delta, S)$  is a commutative, non-cocommutative **Hopf algebra** with **Antipode**  $S(x_1 \cdots x_n) = (-1)^n x_n x_{n-1} \cdots x_1$ .

# Eulerian idempotents

Associative **convolution** in a Hopf Algebra  $(H, \mu, \Delta, S)$ :

$f \star g = \mu \circ (f \otimes g) \circ \Delta$  for any two linear maps  $f, g : H \rightarrow H$ .

For  $H = Sh(A)$  consider **Euler idempotent**  $e : H \rightarrow H$ ,  $e^2 = e$ ,

$$e(x) = \sum_{k=1}^n \frac{(-1)^{k+1}}{k} (\text{id}_H - \epsilon)^{\star k} x \quad \text{for } x \in W_n$$

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**Rem** Closed formula for  $e$  in terms of permutations exists.

**Prop**  $e(x \# y) = 0$  for  $x$  and  $y$  nontrivial.

**Prop** Any word can be decomposed into a unique linear combination of shuffle products of elements from  $\text{ime}$ .

## Eulerian idempotents

**Decomposition** of  $x = x_1 \cdots x_n$

$$x = \sum_{k=1}^n x^{(k)}$$

where

$$x^{(k)} = \frac{1}{k!} \sum_{J_1 \sqcup \cdots \sqcup J_k = \underline{n}} e(x_{J_1}) \# \cdots \# e(x_{J_k})$$

with  $\underline{n} = \{1, \dots, n\}$ ,  $J_i \neq \emptyset$ , and where  $J \sqcup I$  denotes the ordered union (i.e. append the second set to the first) and where for  $J = \{r, r+1, \dots, s\}$ , we set  $x_J = x_r x_{r+1} \cdots x_s$ .

**Grading** from this decomposition:  $\deg(x^{(k)}) = n - k - 1 \geq -1$  for  $x$  a word of length  $n$ . Observe that for any two (nonempty) words  $x, y$  we have  $\deg(x^{(k_1)} \# y^{(k_2)}) = \deg(x^{(k_1)}) + \deg(y^{(k_2)}) + 1$ .



# Poisson structure

**Prop** The image  $\text{ime} \subset Sh(A)$  is a **graded Lie algebra** (w.r.t.  $\text{deg}$ ) with bracket

$$[e(x), e(y)] = \sum_{\rho, \nu \in A} 2\eta^{\rho\nu} e(\partial_\rho^R \star S(x) \cdot S \star \partial_\nu^L(y))$$

where  $\partial_\rho^L(x) = \delta_{\rho, x_1} x_2 \cdots x_n$  and  $\partial_\nu^R(y) = \delta_{\rho, y_m} y_1 \cdots y_{m-1}$

Proof: Antisymmetry easy, Jacobi identity tedious.

$$\begin{aligned} \text{Grading: } [ , ] : e(W_n) \otimes e(W_m) &\rightarrow e(W_{n+m-2}) \\ \Rightarrow \text{deg}([e(x), e(y)]) &= \text{deg}(e(x)) + \text{deg}(e(y)). \end{aligned}$$

Let  $x, y$  be words, decompose them into linear combinations of shuffle products of elements from  $\text{ime}$ , extend the Lie bracket from  $\text{ime}$  via the Leibniz rule  $\rightarrow$  **Poisson bracket on  $Sh(A)$** .

**Prop** On cyclic sums of words this reproduces Poisson bracket of Pohlmeier's  $Z$ .

# Quantization: History

Poisson algebra of  $Z$ 's not freely generated  $\rightarrow$  identify generators and relations. Important tool: grading (work stratum per stratum).

Quantization: deform to a filtered associative s.t. number of relations remains unchanged (admit for quantum corrections)

Never-ending programme (infinitely many strata), but a lot has been learned about the algebra's structure by explicit calculations. Important conjectures were derived.

## Quantization: Breakthrough [Rehren + Meusburger]

again reformulated in  $Sh(A)$ -framework

Let  $V$  be  $d$ -dim vectorspace with basis labelled by  $A$ , i.e.  $v_0, v_1, \dots, v_{d-1}$ . Consider the linear map  $\tilde{\partial} : \bigoplus W_n \rightarrow \bigoplus W_n \oplus V$  (as vectorspaces) given by

$$\tilde{\partial}(x) = (x_2 \cdots x_n, v_{x_1}) - (x_1 \cdots x_{n-1}, v_{x_n}) \quad \text{for } x \in W_n, n \geq 2$$

and  $\tilde{\partial}(x) = 0$  otherwise.

**Prop**  $\tilde{\partial} : \text{ime} \rightarrow \text{ime} \oplus V$  (compatibility with  $e$ )

**Prop**  $\ker \tilde{\partial} = \{\text{cyclically symmetric elements of } \bigoplus W_n\}$ ,  
 i.e.  $\ker \tilde{\partial} = \text{set of physically meaningful objects}$ .

Characterization as kernel of a **derivation**.

Let  $\mathfrak{g} = \text{ime}$  as a Lie algebra.

**Prop**  $V$  becomes a  **$\mathfrak{g}$ -module** via the action

$$\begin{aligned}e(x).v &= 0 \text{ for } x \in W_n, n \neq 2 \text{ for all } v \in V \\e(\mu\rho).v_\nu &= \text{const} (\eta^{\nu\rho} v_\mu - \eta^{\nu\mu} v_\rho) \text{ for } \mu, \rho, \nu \in A\end{aligned}$$

With this action,  $\mathfrak{g} \oplus V$  becomes a **trivial extension**  $\mathfrak{g}_{\text{ext}}$  of  $\mathfrak{g}$ .

*Quantization:* universal enveloping algebras  $U(\mathfrak{g})$  and  $U(\mathfrak{g}_{\text{ext}})$ .

**Prop** The linear map  $\partial : \mathfrak{g} \rightarrow U(\mathfrak{g}_{\text{ext}})$

$$\partial(e(x)) = \frac{1}{2}[v_{x_1}, e(x_2 \cdots x_n)]_+ - \frac{1}{2}[v_{x_n}, e(x_1 \cdots x_{n-1})]_+$$

is a derivation of Lie algebras and can be extended canonically to a derivation  $\partial : U(\mathfrak{g}) \rightarrow U(\mathfrak{g}_{\text{ext}})$ .

**Prop**  $\ker \partial \subset U(\mathfrak{g})$  is an associative subalgebra of  $U(\mathfrak{g})$ .

Let  $z$  be a cyclically symmetric element in  $\oplus W_n$ , i.e.  $z \in \ker \tilde{\partial}$ .  
 consider  $z^{(2)} = \frac{1}{2} \sum_{J \sqcup I = \underline{n}} e(x_J) \# e(x_I)$ .

**Correspondence principle:**

$z^{(2)} \in \ker \partial$  has a uniquely determined *quantum counterpart*  
 $\hat{z}^{(2)} \in \ker \partial$ .

Not in general true for all words in  $\ker \tilde{\partial}$ , but:

Conjecture: Elements of the form  $z^{(2)}$  generate  $\ker \tilde{\partial}$ .

If true, then  $\ker \partial \subset U(\mathfrak{g})$  is a quantum algebra of the **Nambu-Goto string!**. If false, then construction yields quantization of a (large, relevant) subalgebra.

So far, all calculations within the old Pohlmeyer approach reproduced!

# Outlook

Prove conjecture that the set of  $z^{(2)}$  generates  $\ker \tilde{\mathcal{D}}$ .

Even if disproved: prescription yields a quantization of a subalgebra of  $\ker \tilde{\mathcal{D}}$ .

# Outlook

Interesting algebraic questions: classical Poisson algebra is not finitely generated (proved in the 80's), but long standing conjecture that quantum algebra is **finitely generated**. Use the mathematically sound  $\ker \partial \subset U(\mathfrak{g})$  formulation of quantization to check this.

Look for **Hilbert space representation** of  $\ker \partial$ . Role of Casimir operators (which replace the constraints) in the formal closure of the classical algebra of invariants remains to be understood.